

Quadratic relations of the deformed W -algebra for the twisted affine Lie algebra of type $A_{2N}^{(2)}$

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Abstract. We revisit the free field construction of the deformed W -algebra by Frenkel and Reshetikhin, *Commun. Math. Phys.* **197**, 1-31 (1998), where the basic W -current has been identified. Herein, we establish a free field construction of higher W -currents of the deformed W -algebra associated with the twisted affine Lie algebra $A_{2N}^{(2)}$. We obtain a closed set of quadratic relations and duality, which allows us to define deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ using generators and relations.

Key words: deformed W -algebra; twisted affine algebra; quadratic relation; free field construction; exactly solvable model

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1 Introduction

The deformed W -algebra $\mathcal{W}_{x,r}(\mathfrak{g})$ is a two-parameter deformation of the classical W -algebra $\mathcal{W}(\mathfrak{g})$. The deformation theory of the W -algebra has been studied in papers [1, 2, 4, 5, 6, 9, 10, 11, 12, 13, 14, 15]. For instance, free field constructions of the basic W -current $T_1(z)$ of $\mathcal{W}_{x,r}(\mathfrak{g})$ were suggested in the case when the underlying Lie algebra is of classical type. However, in comparison with the conformal case, the deformation theory of W -algebras is still not fully developed and understood. Moreover, finding quadratic relations of the deformed W -algebra $\mathcal{W}_{x,r}(\mathfrak{g})$ is still an unresolved problem.

In this paper, we generalize the study for $\mathcal{W}_{x,r}(A_2^{(2)})$ ¹⁾ by Brazhnikov and Lukyanov [5]. They obtained a quadratic relation for the W -current $T_1(z)$ of the deformed W -algebra $\mathcal{W}_{x,r}(A_2^{(2)})$

$$\begin{aligned} & f\left(\frac{z_2}{z_1}\right) T_1(z_1) T_1(z_2) - f\left(\frac{z_1}{z_2}\right) T_1(z_2) T_1(z_1) \\ &= \delta\left(\frac{x^{-2}z_2}{z_1}\right) T_1(x^{-1}z_2) - \delta\left(\frac{x^2z_2}{z_1}\right) T_1(xz_2) + c \left(\delta\left(\frac{x^{-3}z_2}{z_1}\right) - \delta\left(\frac{x^3z_2}{z_1}\right) \right) \end{aligned}$$

with an appropriate constant c and a function $f(z)$. This study aims to generalize the result for the cases $A_2^{(2)}$ to $A_{2N}^{(2)}$. We introduce higher W -currents $T_i(z)$, $1 \leq i \leq 2N$, by fusion of the free field construction of the basic W -current $T_1(z)$ of $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ [6] (see formula (3.2)). We obtain a closed set of quadratic relations for the W -currents $T_i(z)$, which is completely different from those in the case of deformed W -algebras associated with affine Lie algebras of types $A_N^{(1)}$ and $A(M, N)^{(1)}$ (see formula (3.5)). We refer the reader to references [7, 8] for the affine Lie

¹⁾ We use two types of symbols, $\mathcal{W}_{x,r}(\mathfrak{g})$ and $\mathcal{W}_{x,r}(X_n^{(r)})$, for the deformed W -algebra associated with the affine Lie algebra \mathfrak{g} of type $X_n^{(r)}$.

superalgebra notation. We obtain the duality $T_{2N+1-i}(z) = c_i T_i(z)$ with $1 \leq i \leq N$, which is a new phenomenon that does not occur in the case of deformed W -algebras associated with affine Lie algebras of types $A_2^{(2)}$, $A_N^{(1)}$, and $A(M, N)^{(1)}$ (see formula (3.4)). This allows us to define $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ using generators and relations. We believe that this paper presents a key step toward extending our construction for general affine Lie algebras \mathfrak{g} , because the structures of the free field construction of the basic W -current $T_1(z)$ for the affine algebras other than that of type $A_N^{(1)}$ are quite similar to those of type $A_{2N}^{(2)}$, not $A_N^{(1)}$. We have checked that there are similar quadratic relations as those for type $A_{2N}^{(2)}$ in the case of type $B_N^{(1)}$ with small rank N .

The remainder of this paper is organized as follows. In Section 2, we review the free field construction of the basic W -current $T_1(z)$ of the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ [6]. In Section 3, we introduce higher W -currents $T_i(z)$ and present a closed set of quadratic relations and duality. We also obtain the q -Poisson algebra in the classical limit. In Section 4, we establish proofs of Proposition 3.1 and Theorem 3.2. Section 5 is devoted to discussion. In Appendices A and B, we summarize normal ordering rules.

2 Free field construction

In this section, we define notation and review the free field construction of the basic W -current $T_1(z)$ of $\mathcal{W}_{x,r}(A_{2N}^{(2)})$. Throughout this paper, we fix a natural number $N = 1, 2, 3, \dots$, a real number $r > 1$, and a complex number x with $0 < |x| < 1$.

2.1 Notation

In this section, we use complex numbers a, w, q , and p with $w \neq 0, q \neq 0, \pm 1$, and $|p| < 1$. For any integer n , we define q -integers

$$[n]_q = \frac{q^n - q^{-n}}{q - q^{-1}}.$$

We use symbols for infinite products,

$$(a; p)_\infty = \prod_{k=0}^{\infty} (1 - ap^k), \quad (a_1, a_2, \dots, a_N; p)_\infty = \prod_{i=1}^N (a_i; p)_\infty$$

for complex numbers a_1, a_2, \dots, a_N . The following standard formulae are used,

$$\exp \left(- \sum_{m=1}^{\infty} \frac{1}{m} a^m \right) = 1 - a, \quad \exp \left(- \sum_{m=1}^{\infty} \frac{1}{m} \frac{a^m}{1 - p^m} \right) = (a; p)_\infty.$$

We use the elliptic theta function $\Theta_p(w)$ and the compact notation $\Theta_p(w_1, w_2, \dots, w_N)$,

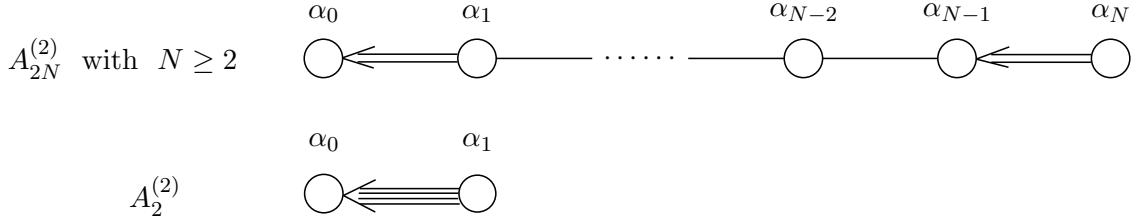
$$\Theta_p(w) = (p, w, pw^{-1}; p)_\infty, \quad \Theta_p(w_1, w_2, \dots, w_N) = \prod_{i=1}^N \Theta_p(w_i)$$

for complex numbers $w_1, w_2, \dots, w_N \neq 0$. Define $\delta(z)$ by the formal series

$$\delta(z) = \sum_{m \in \mathbf{Z}} z^m.$$

2.2 Twisted affine Lie algebra of type $A_{2N}^{(2)}$

In this section we recall the definition of the twisted affine Lie algebra of type $A_{2N}^{(2)}$, $N = 1, 2, 3, \dots$, in Ref.[16]. The Dynkin diagram of type $A_{2N}^{(2)}$ is given by



The corresponding Cartan matrix $A = (A_{i,j})_{i,j=0}^N$ of type $A_{2N}^{(2)}$ is given by

$$A = \begin{pmatrix} 2 & -2 & 0 & \cdots & \cdots & \cdots & 0 \\ -1 & 2 & -1 & \ddots & & & \vdots \\ 0 & -1 & 2 & \ddots & \ddots & & \vdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & & \ddots & \ddots & 2 & -1 & 0 \\ \vdots & & & \ddots & -1 & 2 & -2 \\ 0 & \cdots & \cdots & \cdots & 0 & -1 & 2 \end{pmatrix}$$

with $N \geq 2$, and

$$A = \begin{pmatrix} 2 & -4 \\ -1 & 2 \end{pmatrix}$$

with $N = 1$. We set the labels $a_i = 2$, $0 \leq i \leq N-1$, $a_N = 1$, and the co-labels $a_0^\vee = 1$, $a_i^\vee = 2$, $1 \leq i \leq N$. We set $D = \text{diag}(a_0 a_0^{\vee-1}, a_1 a_1^{\vee-1}, \dots, a_N a_N^{\vee-1})$. We obtain $A = DB$, where B is a symmetric matrix. Thus, the Cartan matrix A is symmetrizable. Let \mathfrak{h} be an $N+2$ -dimensional vector space over \mathbf{C} . Let $\{h_0, h_1, \dots, h_N, d\}$ be a basis of \mathfrak{h} , and $\{\alpha_0, \alpha_1, \dots, \alpha_N, \Lambda_0\}$ a basis of $\mathfrak{h}^* = \text{Hom}_{\mathbf{C}}(\mathfrak{h}, \mathbf{C})$ such that we have with respect to pairing $\langle \cdot, \cdot \rangle : \mathfrak{h} \times \mathfrak{h}^* \rightarrow \mathbf{C}$

$$\langle h_i, \alpha_j \rangle = A_{i,j}, 0 \leq i, j \leq N, \quad \langle d, \alpha_i \rangle = \delta_{0,i}, \quad \langle h_i, \Lambda_0 \rangle = \delta_{i,0}, 0 \leq i \leq N, \quad \langle d, \Lambda_0 \rangle = 0.$$

Let $\mathfrak{g}(A)$ be the affine Lie algebra associated with the Cartan matrix A . Since A is symmetrizable, it is defined as the Lie algebra generated by e_i, f_i , $0 \leq i \leq N$, and \mathfrak{h} with the following relations:

$$\begin{aligned} [e_i, f_j] &= \delta_{i,j} h_i, 0 \leq i, j \leq N, \quad [h, h'] = 0, h, h' \in \mathfrak{h}, \\ [h, e_i] &= \langle h, \alpha_i \rangle e_i, \quad [h, f_i] = -\langle h, \alpha_i \rangle f_i, h \in \mathfrak{h}, 0 \leq i \leq N, \\ (\text{ad } e_i)^{-A_{i,i}+1} e_j &= 0, \quad (\text{ad } f_i)^{-A_{i,i}+1} f_j = 0, 0 \leq i, j \leq N, i \neq j. \end{aligned}$$

Here we used the adjoint action $(\text{ad } x)y = [x, y]$.

2.3 Free field construction

In this section, we recall the free field construction of the basic W -current $T_1(z)$ and of the screening operators S_i of the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ introduced by Frenkel and Reshetikhin [6].

First, we define the $N \times N$ symmetric matrix $B(m) = (B_{i,j}(m))_{i,j=1}^N$, $m \in \mathbf{Z}$, associated with $A_{2N}^{(2)}$, $N = 1, 2, 3, \dots$, as follows.

$$B_{i,j}(m) = \begin{cases} \frac{[2m]_x}{[m]_x}, & 1 \leq i, j \leq N-1, i = j, \\ \frac{[2m]_x - [m]_x}{[m]_x}, & i = j = N, \\ -1, & |i-j| = 1, \\ 0, & |i-j| \geq 2, \end{cases} \quad B_{i,j}(0) = \begin{cases} 2, & 1 \leq i, j \leq N-1, i = j, \\ 1, & i = j = N, \\ -1, & |i-j| = 1, \\ 0, & |i-j| \geq 2. \end{cases}$$

We introduce the Heisenberg algebra $\mathcal{H}_{x,r}$ with generators $a_i(m)$, Q_i , $m \in \mathbf{Z}$, $1 \leq i \leq N$, satisfying

$$\begin{aligned} [a_i(m), a_j(n)] &= \frac{1}{m} [rm]_x [(r-1)m]_x B_{i,j}(m) (x - x^{-1})^2 \delta_{m+n,0}, \quad m, n \neq 0, 1 \leq i, j \leq N, \\ [a_i(0), Q_j] &= B_{i,j}(0), \quad 1 \leq i, j \leq N. \end{aligned}$$

The remaining commutators vanish. The generators $a_i(m)$, Q_i are “root” type generators of $\mathcal{H}_{x,r}$. There is a unique set of “fundamental weight” type generators $y_i(m)$, Q_i^y , $m \in \mathbf{Z}$, $1 \leq i \leq N$, which satisfy the following relations

$$\begin{aligned} [y_i(m), a_j(n)] &= \frac{1}{m} [rm]_x [(r-1)m]_x (x - x^{-1})^2 \delta_{i,j} \delta_{m+n,0}, \quad m, n \neq 0, 1 \leq i, j \leq N, \\ [y_i(0), Q_j^y] &= \delta_{i,j}, \quad [a_i(0), Q_j^y] = \delta_{i,j}, \quad [y_i(0), a_j(m)] = 0, \quad m \in \mathbf{Z}, 1 \leq i, j \leq N. \end{aligned}$$

The explicit formulae for $y_i(m)$ and Q_j^y are given in (A.7). We use the normal ordering : : on $\mathcal{H}_{x,r}$ that satisfies

$$: a_i(m) a_j(n) := \begin{cases} a_i(m) a_j(n), & m < 0, \\ a_j(n) a_i(m), & m \geq 0, \end{cases} \quad m, n \in \mathbf{Z}, 1 \leq i, j \leq N.$$

Let $|0\rangle \neq 0$ be the Fock vacuum of the Fock space of $\mathcal{H}_{x,r}$ such that $a_i(m)|0\rangle = 0$, $m \geq 0$, $1 \leq i \leq N$. Let π_λ be the Fock space of $\mathcal{H}_{x,r}$ generated by $|\lambda\rangle = e^\lambda |0\rangle$, $\lambda = \sum_{j=1}^N \lambda_j Q_j^y$. We obtain

$$a_i(0)|\lambda\rangle = \lambda_i |\lambda\rangle, \quad a_i(m)|\lambda\rangle = 0, \quad m > 0, 1 \leq i \leq N. \quad (2.1)$$

We work in the Fock space π_λ of the Heisenberg algebra $\mathcal{H}_{x,r}$. Let the vertex operators $A_i(z)$, $Y_i(z)$, and $S_i(z)$, $1 \leq i \leq N$, be

$$A_i(z) = x^{ra_i(0)} : \exp \left(\sum_{m \neq 0} a_i(m) z^{-m} \right) :, \quad (2.2)$$

$$Y_i(z) = x^{ry_i(0)} : \exp \left(\sum_{m \neq 0} y_i(m) z^{-m} \right) :, \quad (2.3)$$

$$S_i(z) = z^{\frac{r-1}{2r} B_{i,i}(0)} e^{-\sqrt{\frac{r-1}{r}} Q_i} z^{-\sqrt{\frac{r-1}{r}} a_i(0)} : \exp \left(\sum_{m \neq 0} \frac{a_i(m)}{x^{rm} - x^{-rm}} z^{-m} \right) :. \quad (2.4)$$

The main parts of (2.2), (2.3), and (2.4) are the same as those of Ref.[6]. We corrected the misprints in the formulas for $A_i(z)$, $Y_i(z)$, and $S_i(z)$ in Ref.[6] by multiplying (2.2) and (2.3) by constants and multiplying (2.4) by $z^{\frac{r-1}{2r}B_{i,i}(0)}$. With our fine-tuning, both (3.4) and (3.6) hold.

Let $J_N = \{1, 2, \dots, N, 0, \bar{N}, \dots, \bar{2}, \bar{1}\}$. Here, the indices are ordered as

$$1 \prec 2 \prec \dots \prec N \prec 0 \prec \bar{N} \prec \dots \prec \bar{2} \prec \bar{1}.$$

Let $\bar{k} = k$, $k = 1, 2, \dots, N$, and $\bar{0} = 0$. The indices $i, j \in J_N$ satisfy $i \prec j$ if and only if $\bar{j} \prec \bar{i}$. We define $\bar{I} = \{\bar{i}_1, \bar{i}_2, \dots, \bar{i}_k\}$ for a subset $I \subset J_N$, $I = \{i_1, i_2, \dots, i_k\}$. Let $T_1(z)$ be the generating series with operator valued coefficients acting on the Fock space π_λ ,

$$T_1(z) = \sum_{i \in J_N} \Lambda_i(z), \quad (2.5)$$

where

$$\begin{aligned} \Lambda_1(z) &= Y_1(z), \quad \Lambda_k(z) =: \Lambda_{k-1}(z) A_{k-1}(x^{-k+1}z)^{-1} :, \quad 2 \leq k \leq N, \\ \Lambda_0(z) &= \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x} : \Lambda_N(z) A_N(x^{-N}z)^{-1} :, \\ \Lambda_{\bar{N}}(z) &= \frac{[\frac{1}{2}]_x}{[r - \frac{1}{2}]_x} : \Lambda_0(z) A_N(x^{-N-1}z)^{-1} :, \\ \Lambda_{\bar{k}}(z) &=: \Lambda_{\bar{k}+1}(z) A_k(x^{-2N+k-1}z)^{-1} :, \quad 1 \leq k \leq N-1. \end{aligned} \quad (2.6)$$

We call $T_1(z)$ the basic W -current of the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$.

Let π_μ be the Fock space of $\mathcal{H}_{x,r}$ generated by $|\mu\rangle = e^\mu|0\rangle$ with $\mu = \sum_{i=1}^N \mu_i Q_i^y$, where we choose $\mu_i \in \frac{1}{2}\sqrt{\frac{r-1}{r}}B_{i,i}(0) + \sqrt{\frac{r}{r-1}}\mathbf{Z}$, $1 \leq i \leq N$. From (2.1) and (2.4) the power of w in $S_i(w)$, $w^{\frac{r-1}{2r}B_{i,i}(0)}w^{-\sqrt{\frac{r-1}{r}}a_i(0)}$, takes values in integers on π_μ . Hence, S_i is well-defined on π_μ . We define the screening operators S_i , $1 \leq i \leq N$, acting on the Fock space π_μ as

$$S_i = \oint \frac{dw}{2\pi\sqrt{-1}w} S_i(w). \quad (2.7)$$

The integral in formula (2.7) means the residue at zero.

3 Quadratic relations

In this section, we introduce the higher W -currents $T_i(z)$ and present a set of quadratic relations between $T_i(z)$ for the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$.

3.1 Quadratic relations

We define the formal series $\Delta(z) \in \mathbf{C}[[z]]$ and the constant $c(x, r)$ as

$$\Delta(z) = \frac{(1-x^{2r-1}z)(1-x^{-2r+1}z)}{(1-xz)(1-x^{-1}z)}, \quad c(x, r) = [r]_x[r-1]_x(x-x^{-1}).$$

The formal series $\Delta(z)$ satisfies

$$\begin{aligned} \Delta(z) - \Delta(z^{-1}) &= c(x, r) (\delta(x^{-1}z) - \delta(xz)), \\ \Delta(z)\Delta(x^s z) - \Delta(z^{-1})\Delta(x^{-s} z^{-1}) \\ &= c(x, r) \{ \Delta(x^{s+1})(\delta(x^{-1}z) - \delta(x^{s+1}z)) + \Delta(x^{s-1})(\delta(x^{s-1}z) - \delta(xz)) \}, \quad s \neq 0, \pm 2. \end{aligned}$$

We define the structure functions $f_{i,j}(z)$, $i, j = 0, 1, 2, \dots$, as

$$\begin{aligned} f_{i,j}(z) &= \exp \left(- \sum_{m=1}^{\infty} \frac{1}{m} [(r-1)m]_x [rm]_x (x-x^{-1})^2 \times \right. \\ &\quad \times \left. \frac{[\text{Min}(i,j)m]_x ([N+1-\text{Max}(i,j))m]_x - [(N-\text{Max}(i,j))m]_x}{[m]_x [(N+1)m]_x - [Nm]_x} z^m \right). \end{aligned} \quad (3.1)$$

The ratio of the structure functions $f_{1,1}(z)$ is

$$\frac{f_{1,1}(z^{-1})}{f_{1,1}(z)} = -z \frac{\Theta_{x^{4N+2}}(x^2 z, x^{2N-1} z, x^{4N+2-2r} z, x^{4N+2r} z, x^{2N+1+2r} z, x^{2N-2r+3} z)}{\Theta_{x^{4N+2}}(x^2/z, x^{2N-1}/z, x^{4N+2-2r}/z, x^{4N+2r}/z, x^{2N+1+2r}/z, x^{2N-2r+3}/z)}.$$

We introduce higher W -currents $T_i(z)$ as follows

$$\begin{aligned} T_0(z) &= 1, \quad T_1(z) = \sum_{i \in J_N} \Lambda_i(z), \\ T_i(z) &= \sum_{\substack{\Omega_i \subset J_N \\ |\Omega_i|=i}} d_{\Omega_i}(x, r) \vec{\Lambda}_{\Omega_i}(z), \quad 2 \leq i \leq 2N+1. \end{aligned} \quad (3.2)$$

Here, for a subset $\Omega_i = \{s_1, s_2, \dots, s_i\} \subset J_N$ with $s_1 \prec s_2 \prec \dots \prec s_i$, we set

$$\begin{aligned} d_{\Omega_i}(x, r) &= \prod_{\substack{1 \leq p < q \leq i \\ s_q = s_p}} \Delta(x^{2(q-p+s_p-N-1)}), \quad d_{\emptyset}(x, r) = 1, \\ \vec{\Lambda}_{\Omega_i}(z) &=: \Lambda_{s_1}(x^{-i+1} z) \Lambda_{s_2}(x^{-i+3} z) \cdots \Lambda_{s_i}(x^{i-1} z) :, \quad \vec{\Lambda}_{\emptyset}(z) = 1. \end{aligned} \quad (3.3)$$

Proposition 3.1. *The W -currents $T_i(z)$ satisfy the duality*

$$T_{2N+1-i}(z) = \frac{[r-\frac{1}{2}]_x}{[\frac{1}{2}]_x} \prod_{k=1}^{N-i} \Delta(x^{2k}) T_i(z), \quad 0 \leq i \leq N. \quad (3.4)$$

Theorem 3.2. *The W -currents $T_i(z)$ satisfy the set of quadratic relations*

$$\begin{aligned} &f_{i,j}\left(\frac{z_2}{z_1}\right) T_i(z_1) T_j(z_2) - f_{j,i}\left(\frac{z_1}{z_2}\right) T_j(z_2) T_i(z_1) \\ &= c(x, r) \sum_{k=1}^i \prod_{l=1}^{k-1} \Delta(x^{2l+1}) \left(\delta\left(\frac{x^{-j+i-2k} z_2}{z_1}\right) f_{i-k,j+k}(x^{j-i}) T_{i-k}(x^k z_1) T_{j+k}(x^{-k} z_2) \right. \\ &\quad \left. - \delta\left(\frac{x^{j-i+2k} z_2}{z_1}\right) f_{i-k,j+k}(x^{-j+i}) T_{i-k}(x^{-k} z_1) T_{j+k}(x^k z_2) \right) \\ &+ c(x, r) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) \prod_{l=N+1-j}^{N+i-j} \Delta(x^{2l}) \\ &\times \left(\delta\left(\frac{x^{-2N+j-i-1} z_2}{z_1}\right) T_{j-i}(x^{-i} z_2) - \delta\left(\frac{x^{2N-j+i+1} z_2}{z_1}\right) T_{j-i}(x^i z_2) \right), \quad 1 \leq i \leq j \leq N. \end{aligned} \quad (3.5)$$

Here, we use $f_{i,j}(z)$ introduced in (3.1).

In view of Proposition 3.1 and Theorem 3.2, we obtain the following definition.

Definition 3.3. Let W be the free complex associative algebra generated by elements $\bar{T}_i[m], m \in \mathbf{Z}, 1 \leq i \leq 2N$, I_K the left ideal generated by elements $\bar{T}_i[m], m \geq K \in \mathbf{N}, 1 \leq i \leq 2N$, and

$$\widehat{W} = \lim_{\leftarrow} W/I_K.$$

The deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ is the quotient of \widehat{W} by the two-sided ideal generated by the coefficients of the generating series which are the differences of the right hand sides and of the left hand sides of the relations (3.4) and (3.5), where the generating series $T_i(z)$ are replaced with $\bar{T}_i(z) = \sum_{m \in \mathbf{Z}} \bar{T}_i[m]z^{-m}, 1 \leq i \leq 2N$, and $\bar{T}_0(z) = 1$.

The justification of this definition is presented later. We compare this definition of the deformed W -algebra with other definitions in Section 5.

Lemma 3.4. *The W -currents $T_i(z)$ commute with the screening operators S_j ,*

$$[T_i(z), S_j] = 0, \quad 1 \leq i \leq 2N, 1 \leq j \leq N. \quad (3.6)$$

We present the proofs of Proposition 3.1, Theorem 3.2, and Lemma 3.4 in Section 4.

3.2 Classical limit

The deformed W -algebra $\mathcal{W}_{x,r}(\mathfrak{g})$ yields a q -Poisson W -algebra [6, 17, 18, 19] in the classical limit. As an application of the quadratic relations (3.5), we obtain a q -Poisson W -algebra of type $A_{2N}^{(2)}$. We set parameters $q = x^{2r}$ and $\beta = (r - 1)/r$. We define the q -Poisson bracket $\{\cdot, \cdot\}$ by taking the classical limit $\beta \rightarrow 0$ with q fixed as

$$\{T_i^{\text{PB}}[m], T_j^{\text{PB}}[n]\} = \lim_{\beta \rightarrow 0} \frac{1}{2\beta \log q} [T_i[m], T_j[n]].$$

Here, we introduce $T_i^{\text{PB}}[m]$ by $T_i(z) = \sum_{m \in \mathbf{Z}} T_i[m]z^{-m} \longrightarrow T_i^{\text{PB}}(z) = \sum_{m \in \mathbf{Z}} T_i^{\text{PB}}[m]z^{-m}, \beta \rightarrow 0, q \text{ fixed.}$

The β -expansions of the structure functions are given as

$$\begin{aligned} f_{i,j}(z) &= 1 - 2\beta \log q (q - q^{-1}) \sum_{m=1}^{\infty} [\text{Min}(i, j)m]_q \times \\ &\quad \times \frac{[(N+1-\text{Max}(i, j))m]_q - [(N-\text{Max}(i, j))m]_q}{[(N+1)m]_q - [Nm]_q} z^m + O(\beta^2), \quad i, j \geq 1, \\ c(x, r) &= 2\beta \log q + O(\beta^2). \end{aligned}$$

As corollaries of Proposition 3.1 and Theorem 3.2 we obtain the following.

Corollary 3.5. *For the q -Poisson W -algebra associated with affine Lie algebra of type $A_{2N}^{(2)}$, the currents $T_i^{\text{PB}}(z)$ satisfy*

$$\begin{aligned} \{T_i^{\text{PB}}(z_1), T_j^{\text{PB}}(z_2)\} &= (q - q^{-1}) C_{i,j} \left(\frac{z_2}{z_1} \right) T_i^{\text{PB}}(z_1) T_j^{\text{PB}}(z_2) \\ &\quad + \sum_{k=1}^i \left(\delta \left(\frac{q^{-j+i-2k} z_2}{z_1} \right) T_{i-k}^{\text{PB}}(q^k z_1) T_{j+k}^{\text{PB}}(q^{-k} z_2) - \delta \left(\frac{q^{j-i+2k} z_2}{z_1} \right) T_{i-k}^{\text{PB}}(q^{-k} z_1) T_{j+k}^{\text{PB}}(q^k z_2) \right) \\ &\quad + \delta \left(\frac{q^{-2N+j-i-1} z_2}{z_1} \right) T_{j-i}^{\text{PB}}(q^{-i} z_2) - \delta \left(\frac{q^{2N-j+i+1} z_2}{z_1} \right) T_{j-i}^{\text{PB}}(q^i z_2), \quad 1 \leq i \leq j \leq N. \quad (3.7) \end{aligned}$$

Here, the structure functions $C_{i,j}(z)$ are given by

$$C_{i,j}(z) = \sum_{m \in \mathbf{Z}} \frac{[\text{Min}(i, j)m]_q ((N+1-\text{Max}(i, j))m]_q - [(N-\text{Max}(i, j))m]_q)}{[(N+1)m]_q - [Nm]_q} z^m, \quad 1 \leq i, j \leq N.$$

Corollary 3.6. *The currents $T_i^{PB}(z)$ satisfy the duality relations*

$$T_{2N+1-i}^{PB}(z) = T_i^{PB}(z), \quad 0 \leq i \leq N. \quad (3.8)$$

4 Proof of Theorem 3.2

In this section, we prove Proposition 3.1, Theorem 3.2, and Lemma 3.4.

4.1 Proof of Proposition 3.1

Lemma 4.1. *The $\Lambda_i(z)$, $i \in J_N$, satisfy*

$$\begin{aligned} f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_i(z_1) \Lambda_j(z_2) &= \Delta \left(\frac{x^{-1} z_2}{z_1} \right) : \Lambda_i(z_1) \Lambda_j(z_2) :, \quad i, j \in J_N, i \prec j, j \neq \bar{i}, \\ f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_j(z_1) \Lambda_i(z_2) &= \Delta \left(\frac{x z_2}{z_1} \right) : \Lambda_j(z_1) \Lambda_i(z_2) :, \quad i, j \in J_N, i \prec j, j \neq \bar{i}, \\ f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_0(z_1) \Lambda_0(z_2) &= \Delta \left(\frac{z_2}{z_1} \right) : \Lambda_0(z_1) \Lambda_0(z_2) :, \\ f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_i(z_1) \Lambda_i(z_2) &=: \Lambda_i(z_1) \Lambda_i(z_2) :, \quad i \in J_N \setminus \{0\}, \\ f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_k(z_1) \Lambda_{\bar{k}}(z_2) &= \Delta \left(\frac{x^{-1} z_2}{z_1} \right) \Delta \left(\frac{x^{-2N-2+2k} z_2}{z_1} \right) : \Lambda_k(z_1) \Lambda_{\bar{k}}(z_2) :, \quad 1 \leq k \leq N, \\ f_{1,1} \left(\frac{z_2}{z_1} \right) \Lambda_{\bar{k}}(z_1) \Lambda_k(z_2) &= \Delta \left(\frac{x z_2}{z_1} \right) \Delta \left(\frac{x^{2N+2-2k} z_2}{z_1} \right) : \Lambda_{\bar{k}}(z_1) \Lambda_k(z_2) :, \quad 1 \leq k \leq N. \end{aligned} \quad (4.1)$$

Proof. Using (A.2) and (A.8), we obtain the normal ordering rules (4.1). ■

Lemma 4.2. *The $\Lambda_i(z)$, $i \in J_N$, satisfy*

$$: \Lambda_0(z) \Lambda_0(xz) := \Delta(1) : \Lambda_N(z) \Lambda_{\bar{N}}(xz) :, \quad (4.2)$$

$$: \Lambda_1(z) \Lambda_{\bar{1}}(x^{2N+1} z) := 1, \quad (4.3)$$

$$: \Lambda_k(z) \Lambda_{\bar{k}}(x^{2N-2k+3} z) :=: \Lambda_{k-1}(z) \Lambda_{\bar{k}-1}(x^{2N-2k+3} z) :, \quad 2 \leq k \leq N. \quad (4.4)$$

Proof. From (2.6), we obtain (4.2) and (4.4). From (2.2), (2.3) and (2.6), we obtain (4.3). ■

Lemma 4.3. *The $\Delta(z)$ and $f_{i,j}(z)$ satisfy the following fusion relations.*

$$f_{i,j}(z) = f_{j,i}(z) = \prod_{k=1}^i f_{1,j}(z^{-i-1+2k} z), \quad 1 \leq i \leq j, \quad (4.5)$$

$$f_{1,i}(z) = \left(\prod_{k=1}^{i-1} \Delta(x^{-i+2k} z) \right)^{-1} \prod_{k=1}^i f_{1,1}(x^{-i-1+2k} z), \quad i \geq 2, \quad (4.6)$$

$$f_{i,2N+1}(z) = \prod_{k=1}^i \Delta(x^{-i-1+2k} z), \quad i \geq 1, \quad (4.7)$$

$$f_{i,j}(z) = f_{i,2N+1-j}(z) = f_{2N+1-j,i}(z) = f_{j,i}(z), \quad i \geq 1, 1 \leq j \leq N, \quad (4.8)$$

$$f_{1,j}(z) f_{1,j}(x^{2N+1} z) = \Delta(x^j z) \Delta(x^{2N+1-j} z), \quad j \geq 1, \quad (4.9)$$

$$f_{1,i}(z) f_{j,i}(x^{\pm(j+1)} z) = \begin{cases} f_{j+1,i}(x^{\pm j} z) \Delta(x^{\pm i} z), & 1 \leq i \leq j, \\ f_{j+1,i}(x^{\pm j} z), & 1 \leq j < i, \end{cases} \quad (4.10)$$

$$f_{1,i}(z) f_{1,j}(x^{\pm(i+j)} z) = f_{1,i+j}(x^{\pm j} z) \Delta(x^{\pm i} z), \quad i, j \geq 1, \quad (4.11)$$

$$f_{1,i}(z) f_{1,j}(x^{\pm(i-j-2k)} z) = f_{1,i-k}(x^{\mp k} z) f_{1,j+k}(x^{\pm(i-j-k)} z), \quad i, j, i-k, j+k \geq 1. \quad (4.12)$$

Proof. We show (4.6) here. From the definitions, we have

$$\begin{aligned} & \left(\prod_{k=1}^{i-1} \Delta_1(x^{-i+2k}z) \right)^{-1} \prod_{k=1}^i f_{1,1}(x^{-i-1+2k}z) \\ &= \exp \left(- \sum_{m=1}^{\infty} \frac{1}{m} \frac{[rm]_x[(r-1)m]_x}{[(N+1)m]_x - [Nm]_x} (x - x^{-1})^2 \times \right. \\ & \quad \left. \times \left\{ ([Nm]_x - [(N-1)m]_x) \sum_{k=1}^i x^{(-i+2k-1)m} - ([(N+1)m]_x - [Nm]_x) \sum_{k=1}^{i-1} x^{(-i+2k)m} \right\} z^m \right). \end{aligned}$$

Using the relation

$$[(a-1)m]_x \sum_{k=1}^i x^{(-i+2k-1)m} - [am]_x \sum_{k=1}^{i-1} x^{(-i+2k)m} = [(a-i)m]_x, \quad a = N, N+1,$$

we obtain $f_{1,i}(z)$ in the right hand side of the previous formula. We obtain (4.5), (4.7), (4.8), and (4.9) by straightforward calculation from the definitions. Using (4.5) and (4.6), we obtain the relations (4.10), (4.11), and (4.12). \blacksquare

Lemma 4.4. *The following relation holds for $A \subset J_N$,*

$$\overrightarrow{\Lambda}_{\overline{J_N \setminus A}}(z) = \overrightarrow{\Lambda}_A(z) \times \begin{cases} \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x}, & 0 \notin A, \\ \frac{[\frac{1}{2}]_x}{[r - \frac{1}{2}]_x}, & 0 \in A. \end{cases} \quad (4.13)$$

Proof. First, we consider the case $A = \emptyset$ and $\overline{J_N \setminus A} = J_N$. In this case, (4.13) can be rewritten as

$$: \Lambda_1(x^{-2N}z) \cdots \Lambda_N(x^{-2}z) \Lambda_0(z) \Lambda_{\overline{N}}(x^2z) \cdots \Lambda_{\overline{1}}(x^{2N}z) : = \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x}. \quad (4.14)$$

Using (2.2), (2.3), and (2.6), the left side of (4.14) can be written as

$$\frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x} : \exp \left(\sum_{m \neq 0} \left(\frac{[(2N+1)m]_x}{[m]_x} y_1(m) - \sum_{j=1}^N \frac{[(2N+1-j)m]_x + [jm]_x}{[m]_x} a_j(m) \right) z^{-m} \right) : .$$

Using the relation $\frac{[(2N+1-j)m]_x + [jm]_x}{[(2N+1)m]_x} = \frac{[(N+1-j)m]_x - [(N-j)m]_x}{[(N+1)m]_x - [Nm]_x}$, the generators $y_1(m)$ in (A.7) are rewritten as $y_1(m) = \sum_{j=1}^N \frac{[(2N+1-j)m]_x + [jm]_x}{[(2N+1)m]_x} a_j(m)$. Hence, we obtain (4.14).

Next, we show (4.13) for $A \subset J_N$. Cases (i), $0 \in A$ and (ii), $0 \notin A$ are proved separately. First, we study case (i), $0 \in A$. Let

$$A = \{k_1, \dots, k_K, 0, \overline{l_L}, \dots, \overline{l_1} \mid k_1 \prec \dots \prec k_K \prec 0 \prec \overline{l_L} \prec \dots \prec \overline{l_1}, \quad 1 \leq K, L \leq N\}.$$

Multiplying (4.14) by $\overrightarrow{\Lambda}_A(x^{L-K+1}z)$ on the left, and using (4.1) and (4.7) yields

$$: \overrightarrow{\Lambda}_{J_N}(z) \overrightarrow{\Lambda}_A(x^{L-K+1}z) : = \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x} \overrightarrow{\Lambda}_A(x^{L-K+1}z). \quad (4.15)$$

Using (4.2), (4.3) and (4.4) yields

$$\begin{aligned}
& : \overrightarrow{\Lambda}_{J_N}(z)\Lambda_0(xz) := \Delta(1)\overrightarrow{\Lambda}_{J_N \setminus \{0\}}(xz), \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{0\}}(z)\Lambda_{\bar{l}_L}(x^2 z) := \overrightarrow{\Lambda}_{J_N \setminus \{l_L, 0\}}(xz), \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_{L-s+1}, \dots, l_L, 0\}}(z)\Lambda_{\bar{l}_{L-s}}(x^{2+s} z) := \overrightarrow{\Lambda}_{J_N \setminus \{l_{L-s}, \dots, l_L, 0\}}(xz), \quad 1 \leq s \leq L-1, \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, 0\}}(z)\Lambda_{k_K}(x^{-L-2} z) := \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, 0, \bar{k}_K\}}(x^{-1} z), \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, 0, \bar{k}_K, \dots, \bar{k}_{K-s+1}\}}(z)\Lambda_{k_{K-s}}(x^{-L-2-s} z) : \\
& = \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, 0, \bar{k}_K, \dots, \bar{k}_{K-s}\}}(x^{-1} z), \quad 1 \leq s \leq K-1.
\end{aligned}$$

Using the above five relations yields

$$: \overrightarrow{\Lambda}_{J_N}(z)\overrightarrow{\Lambda}_A(x^{L-K+1} z) := \Delta(1)\overrightarrow{\Lambda}_{\overline{J_N \setminus A}}(x^{L-K+1} z).$$

From (4.15) we obtain (4.13) for $0 \in A$.

Next, we study case (ii), $0 \notin A$. The proof for this case is similar to that of case (i). Let

$$A = \{k_1, \dots, k_K, \bar{l}_L, \dots, \bar{l}_1 \mid k_1 \prec \dots \prec k_K \prec \bar{l}_L \prec \dots \prec \bar{l}_1, \quad 1 \leq K, L \leq N\}.$$

Multiplying (4.14) by $\overrightarrow{\Lambda}_A(x^{L-K} z)$ on the left, and using (4.1) and (4.7) yields

$$: \overrightarrow{\Lambda}_{J_N}(z)\overrightarrow{\Lambda}_A(x^{L-K} z) := \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x} \overrightarrow{\Lambda}_A(x^{L-K} z). \quad (4.16)$$

Using (4.2), (4.3), and (4.4) yields

$$\begin{aligned}
& : \overrightarrow{\Lambda}_{J_N}(z)\Lambda_{\bar{l}_L}(xz) := \overrightarrow{\Lambda}_{J_N \setminus \{l_L\}}(xz), \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_{L-s+1}, \dots, l_L\}}(z)\Lambda_{\bar{l}_{L-s}}(x^{1+s} z) := \overrightarrow{\Lambda}_{J_N \setminus \{l_{L-s}, \dots, l_L\}}(xz), \quad 1 \leq s \leq L-1, \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L\}}(z)\Lambda_{k_K}(x^{-L-1} z) := \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, \bar{k}_K\}}(x^{-1} z), \\
& : \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, \bar{k}_K, \dots, \bar{k}_{K-s+1}\}}(z)\Lambda_{k_{K-s}}(x^{-L-1-s} z) : \\
& = \overrightarrow{\Lambda}_{J_N \setminus \{l_1, \dots, l_L, \bar{k}_K, \dots, \bar{k}_{K-s}\}}(x^{-1} z), \quad 1 \leq s \leq K-1.
\end{aligned}$$

Using the above five relations yields

$$: \overrightarrow{\Lambda}_{J_N}(z)\overrightarrow{\Lambda}_A(x^{L-K} z) := \overrightarrow{\Lambda}_{\overline{J_N \setminus A}}(x^{L-K} z).$$

From (4.16) we obtain (4.13) for $0 \notin A$. ■

Lemma 4.5. *The following relation holds for $A \subset J_N$ with $|A| \leq N$:*

$$\frac{d_{J_N \setminus A}(x, r)}{d_A(x, r)} = \prod_{k=1}^{N-|A|} \Delta(x^{2k}) \times \begin{cases} \Delta(1), & 0 \in A, \\ 1, & 0 \notin A. \end{cases} \quad (4.17)$$

Proof. We define the map $\sigma : J_N \rightarrow J_{N+1}$ by

$$\sigma(j) = \begin{cases} k+1, & j = k, 1 \leq k \leq N, \\ 0, & j = 0, \\ \bar{k}+1, & j = \bar{k}, 1 \leq k \leq N. \end{cases}$$

For $T \subset J_N$ with $|T| \leq N$, relation (4.17) is rewritten as

$$\frac{d_{\sigma(J_N \setminus T)}(x, r)}{d_{\sigma(T)}(x, r)} = \prod_{k=1}^{N-|T|} \Delta(x^{2k}) \times \begin{cases} \Delta(1), & 0 \in T, \\ 1, & 0 \notin T. \end{cases}$$

Hence, the relation

$$\frac{d_{(J_N \setminus B) \cap \sigma(J_N)}(x, r)}{d_{B \cap \sigma(J_N)}(x, r)} = \prod_{k=1}^{N-|B \cap \sigma(J_N)|} \Delta(x^{2k}) \times \begin{cases} \Delta(1), & 0 \in B, \\ 1, & 0 \notin B, \end{cases} \quad (4.18)$$

for $B \subset J_{N+1}$ with $|B \cap \sigma(J_N)| \leq N$ holds if relation (4.17) for $A \subset J_N$ with $|A| \leq N$ is assumed. Here, we used $|B \cap \sigma(J_N)| = |\sigma^{-1}(B \cap \sigma(J_N))|$.

We prove (4.17) by induction on N . First, we establish the base $N = 1$ using case-by-case analysis. For $A = \emptyset$, we obtain $d_A(x, r) = 1$ and $d_{J_1 \setminus A}(x, r) = \Delta(x^2)$. For $A = \{1\}$, we obtain $d_A(x, r) = 1$ and $d_{J_1 \setminus A}(x, r) = 1$. For $A = \{0\}$, we obtain $d_A(x, r) = 1$ and $d_{J_1 \setminus A}(x, r) = \Delta(1)$. For $A = \{\bar{1}\}$, we obtain $d_A(x, r) = 1$ and $d_{J_1 \setminus A}(x, r) = 1$. This implies that (4.17) holds for $N = 1$.

Next, we assume that relation (4.17) holds for some N , and show (4.17) for N replaced by $N + 1$. Let $A \subset J_{N+1}$. From the definition of $d_A(x, r)$, we obtain

$$\frac{d_{J_N \setminus A}(x, r)}{d_A(x, r)} = \frac{d_{(J_N \setminus A) \cap \sigma(J_N)}(x, r)}{d_{A \cap \sigma(J_N)}(x, r)} \times \begin{cases} 1, & 1 \in A, \bar{1} \notin A \text{ or } 1 \notin A, \bar{1} \in A, \\ \Delta(x^{2(N-|J_N \setminus A|+1)})^{-1}, & 1, \bar{1} \in A, \\ \Delta(x^{2(N-|A|+1)}), & 1, \bar{1} \notin A. \end{cases} \quad (4.19)$$

Cases (i), $1 \in A, \bar{1} \notin A$ (or $1 \notin A, \bar{1} \in A$), (ii), $1, \bar{1} \in A$, and (iii), $1, \bar{1} \notin A$ are proved separately.

First, we study case (i), $1 \in A, \bar{1} \notin A$ (or $1 \notin A, \bar{1} \in A$). In this case, we obtain $|A \cap \sigma(J_N)| = |A| - 1 \leq N$. Hence, (4.18) holds with $B = A$. Using (4.18), (4.19) and $|A \cap \sigma(J_N)| = |A| - 1$ yields (4.17) with N replaced by $N + 1$.

Next, we study case (ii), $1, \bar{1} \in A$. In this case, we obtain $|A \cap \sigma(J_N)| = |A| - 2 \leq N - 1$. Hence, (4.18) holds with $B = A$. Using (4.18) and (4.19), $|A \cap \sigma(J_N)| = |A| - 2$ and $|J_N \setminus A| = 2N + 3 - |A|$ yields (4.17) with N replaced by $N + 1$.

Finally, we examine case (iii), $1, \bar{1} \notin A$. Case (iii) is further subdivided into (iii.1), $|A| \leq N$, $1, \bar{1} \notin A$ and (iii.2), $|A| = N + 1$, $1, \bar{1} \notin A$.

For the condition (iii.1), we obtain $|A \cap \sigma(J_N)| = |A| \leq N$. Hence, (4.18) holds with $B = A$. Using (4.18), (4.19), and $|A \cap \sigma(J_N)| = |A|$ yields (4.17) with N replaced by $N + 1$.

For condition (iii.2), we obtain $|(J_N \setminus A) \cap \sigma(J_N)| = N$. Hence, (4.18) holds with $B = J_N \setminus A$. Using (4.18) and (4.19), $|A| = N + 1$ and $|(J_N \setminus A) \cap \sigma(J_N)| = N$ yields (4.17) with N replaced by $N + 1$. \blacksquare

Proof. Here we will show Proposition 3.1. Using (4.13), (4.17), and $d_{\overline{J_N \setminus \Omega_i}}(x, r) = d_{J_N \setminus \Omega_i}(x, r)$ yields

$$d_{\overline{J_N \setminus \Omega_i}}(x, r) \overrightarrow{\Lambda}_{\overline{J_N \setminus \Omega_i}}(z) = \frac{[r - \frac{1}{2}]_x}{[\frac{1}{2}]_x} \prod_{k=1}^{N-|\Omega_i|} \Delta(x^{2k}) d_{\Omega_i}(x, r) \overrightarrow{\Lambda}_{\Omega_i}(z). \quad (4.20)$$

Adding relations (4.20) over all $\Omega_i \subset J_N$ for each fixed i , $0 \leq i \leq N$, yields (3.4). \blacksquare

4.2 Proof of Theorem 3.2

Lemma 4.6. *The W -currents $T_j(z)$, $1 \leq j \leq N$, satisfy the set of quadratic relations*

$$\begin{aligned} & f_{1,j} \left(\frac{z_2}{z_1} \right) T_1(z_1) T_j(z_2) - f_{j,1} \left(\frac{z_1}{z_2} \right) T_j(z_2) T_1(z_1) \\ &= c(x, r) \left(\delta \left(\frac{x^{-j-1} z_2}{z_1} \right) T_{j+1}(x^{-1} z_2) - \delta \left(\frac{x^{j+1} z_2}{z_1} \right) T_{j+1}(x z_2) \right) \\ &+ c(x, r) \Delta(x^{2N+2-2j}) \left(\delta \left(\frac{x^{-2N+j-2} z_2}{z_1} \right) T_{j-1}(x^{-1} z_2) - \delta \left(\frac{x^{2N-j+2} z_2}{z_1} \right) T_{j-1}(x z_2) \right), \\ & 1 \leq j \leq N. \end{aligned} \quad (4.21)$$

Here, we use $f_{i,j}(z)$ introduced in (3.1).

Proof. In this proof, we frequently use exchange relations (B.2)–(B.8) in Appendix B. We start from

$$\text{LHS}_{1,j} = f_{1,j}(z_2/z_1) T_1(z_1) T_j(z_2) - f_{j,1}(z_1/z_2) T_j(z_2) T_1(z_1), \quad 1 \leq j \leq N. \quad (4.22)$$

From the definition of $T_j(z)$ introduced in (3.2), $\text{LHS}_{1,j}$ can be written as the sum of

$$f_{1,j}(z_2/z_1) \Lambda_s(z_1) \vec{\Lambda}_{\Omega_j}(z_2) - f_{j,1}(z_1/z_2) \vec{\Lambda}_{\Omega_j}(z_2) \Lambda_s(z_1) \text{ over } s \in J_N, \Omega_j \subset J_N, |\Omega_j| = j \quad (4.23)$$

summarized in Appendix B. Adding exchange relations (B.2)–(B.8) over $s \in J_N, \Omega_j \subset J_N, |\Omega_j| = j$ yields

$$\begin{aligned} \text{LHS}_{1,j} &= c(x, r) \left\{ \sum_{m=0}^{\lfloor \frac{j}{2} \rfloor} \left(\delta \left(x^{-j-1+2m} \frac{z_2}{z_1} \right) \bar{G}_{j+1-2m}(z_2) - \delta \left(x^{j+1-2m} \frac{z_2}{z_1} \right) G_{j+1-2m}(z_2) \right) \right. \\ &+ \left. \sum_{m=0}^{N-\lfloor \frac{j-1}{2} \rfloor} \left(\delta \left(x^{-2N+j-2+2m} \frac{z_2}{z_1} \right) \bar{H}_{2N-j+2-2m}(z_2) - \delta \left(x^{2N-j+2-2m} \frac{z_2}{z_1} \right) H_{2N-j+2-2m}(z_2) \right) \right\}. \end{aligned} \quad (4.24)$$

Formulas for $\bar{G}_{j+1}(z)$, $G_{j+1}(z)$, $\bar{H}_{2N-j+2}(z)$, $H_{2N-j+2}(z)$, $\bar{G}_{j+1-2m}(z)$, $G_{j+1-2m}(z)$, $\bar{H}_{2N-j+2-2m}(z)$, and $H_{2N-j+2-2m}(z)$ will be given below. In (4.45) we define $H_0(z) = 0$ to avoid ambiguity of $\bar{H}_0(z)$ and $H_0(z)$. In the case when j is even, we have $\text{LHS}_{1,j} = c(x, r) (\bar{H}_0(z_2) - H_0(z_2)) \delta(z_2/z_1) + \bar{H}_2(z_2) \delta(x^{-2} z_2/z_1) - H_2(z_2) \delta(x^2 z_2/z_1) + \dots$.

First, we define $\bar{G}_{j+1}(z)$, $1 \leq j \leq N$, as the coefficient of $\delta(x^{-j-1} z_2/z_1)$ in (4.24). In what follows, for a subset $\Omega_j \subset J_N$ with $|\Omega_j| = j$, we write its elements as s_1, s_2, \dots, s_j , $s_1 \prec s_2 \prec \dots \prec s_j$. Adding the first term in (B.2) and the first term in (B.7) yields

$$\begin{aligned} \bar{G}_{j+1}(z) &= \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s \prec s_1, \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1} z) \vec{\Lambda}_{\Omega_j}(z) : \\ &+ \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{l=1 \\ n \prec s_1 \\ s_l=\bar{n}}}^j \Delta(x^{2(N+1-l-n)}) d_{\Omega_j}(x, r) : \Lambda_n(x^{-j-1} z) \vec{\Lambda}_{\Omega_j}(z) :. \end{aligned} \quad (4.25)$$

Using $: \Lambda_n(x^{-j-1} z) \vec{\Lambda}_{\Omega_j}(z) := \vec{\Lambda}_{\Omega_j \cup \{n\}}(x^{-1} z)$ and

$$d_{\Omega_j \cup \{n\}}(x, r) = d_{\Omega_j}(x, r) \times \begin{cases} 1, & \bar{n} \notin \Omega_j, \\ \Delta(x^{2(N+1-l-n)}), & \bar{n} = s_l \end{cases} \quad \text{with } n \prec s_1, 1 \leq n \leq N,$$

yields

$$\overline{G}_{j+1}(z) = \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s \prec s_1}} d_{\Omega_j \cup \{s\}}(x, r) \overrightarrow{\Lambda}_{\Omega_j \cup \{s\}}(x^{-1}z).$$

Hence, we obtain $\overline{G}_{j+1}(z) = T_{j+1}(x^{-1}z), 1 \leq j \leq N$.

Next, we define $G_{j+1}(z), 1 \leq j \leq N$, as the coefficient of $\delta(x^{j+1}z_2/z_1)$ in (4.24). Adding the second term in (B.2) and the third term in (B.8) yields

$$\begin{aligned} G_{j+1}(z) &= \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_j \prec s, \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1}z) : \\ &\quad + \sum_{\Omega_j \subset J_N} \sum_{\substack{n=1 \\ s_j \prec \bar{n}}}^N \sum_{k=1}^j \Delta(x^{2(N+k-j-n)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{j+1}z) : . \end{aligned} \quad (4.26)$$

Using : $\overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{j+1}z) := \overrightarrow{\Lambda}_{\Omega_j \cup \{\bar{n}\}}(xz)$ and

$$d_{\Omega_j \cup \{\bar{n}\}}(x, r) = d_{\Omega_j}(x, r) \times \begin{cases} 1, & n \notin \Omega_j, \\ \Delta(x^{2(N+k-j-n)}), & n = s_k \end{cases} \quad \text{with } s_j \prec \bar{n}, 1 \leq n \leq N,$$

yields

$$G_{j+1}(z) = \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_j \prec s}} d_{\Omega_j \cup \{s\}}(x, r) \overrightarrow{\Lambda}_{\Omega_j \cup \{s\}}(xz).$$

Hence, we obtain $G_{j+1}(z) = T_{j+1}(z), 1 \leq j \leq N$.

We define $\overline{H}_{2N-j+2}(z), 1 \leq j \leq N$, as the coefficient of $\delta(x^{-2N+j-2}z_2/z_1)$ in (4.24). Adding the first term in (B.4), the second term in (B.4), the second term in (B.7), and the fourth term in (B.7) yields

$$\begin{aligned} \overline{H}_{2N-j+2}(z) &= \sum_{n=1}^{j-1} \sum_{k=1}^{j-n} \sum_{\substack{\Omega_j \subset J_N \\ s_k = n \\ s_l = \bar{n}, l=j-n+1}} d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2}z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\ &\quad - \sum_{n=1}^{j-2} \sum_{k=1}^{j-n-1} \sum_{\substack{\Omega_j \subset J_N \\ s_k = n \\ s_l = \bar{n}, l=j-n}} d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2}z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\ &\quad + \sum_{n=1}^j \sum_{k=1}^{j-n+1} \sum_{\substack{\Omega_j \subset J_N \\ s_{k-1} \prec n \prec s_k \\ s_l = \bar{n}, l=j+1-n}} \Delta(x^{2(N-j+k)}) d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2}z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\ &\quad - \sum_{n=1}^{j-1} \sum_{k=1}^{j-n} \sum_{\substack{\Omega_j \subset J_N \\ s_{k-1} \prec n \prec s_k \\ s_l = \bar{n}, l=j-n}} \Delta(x^{2(N-j+k)}) d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2}z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \end{aligned} \quad (4.27)$$

The second term in (4.27) vanishes, because there doesn't exist $s_j \in J_N$, if $\Omega_j \subset J_N, s_k = n, 1 \leq k \leq j+1-n, s_l = \bar{n}, l = j-n$, and $1 \leq n \leq j-2$ are satisfied. The fourth term in (4.27) vanishes,

because there doesn't exist $s_j \in J_N$, if $\Omega_j \subset J_N$, $s_{k-1} \prec n \prec s_k$, $1 \leq k \leq j-n$, $s_l = \bar{n}$, $l = j-n$, and $1 \leq n \leq j-1$ are satisfied. Rewriting the sum of the first and the third terms yields

$$\begin{aligned} \overline{H}_{2N-j+2}(z) &= \sum_{n=1}^{j-1} \sum_{k=1}^{\text{Min}(j-n, n)} \sum_{\substack{\Omega_j \subset J_N \\ s_k=n \\ (s_{j-n+1}, \dots, s_{j-1}, s_j) = (\bar{n}, \dots, \bar{2}, \bar{1})}} d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\ &\quad + \sum_{n=1}^j \sum_{k=1}^{\text{Min}(j+1-n, n)} \sum_{\substack{\Omega_j \subset J_N \\ s_{k-1} \prec n \prec s_k \\ (s_{j-n+1}, \dots, s_{j-1}, s_j) = (\bar{n}, \dots, \bar{2}, \bar{1})}} \Delta(x^{2(N-j+k)}) d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \end{aligned}$$

The relation $d_{\Omega_j}(x, r) = \Delta(x^{2(N+1-j)}) d_{\Omega_j \setminus \{\bar{n}\}}(x, r)$ holds, if $s_k = n$, $s_l = \bar{n}$, $1 \leq n \leq j-1$, and $1 \leq k < l = j+1-n$ are satisfied. The relation $d_{\Omega_j}(x, r) \Delta(x^{2(N-j+k)}) = \Delta(x^{2(N+1-j)}) d_{\Omega_j \setminus \{\bar{n}\}}(x, r)$ holds, if $s_{k-1} \prec n \prec s_k$, $s_l = \bar{n}$, $1 \leq n \leq j$, and $1 \leq k \leq l = j+1-n$ are satisfied. Using the above two relations and

$$: \Lambda_n(x^{-2N+j-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) := \overrightarrow{\Lambda}_{\Omega_j \setminus \{\bar{n}\}}(x^{-1} z), \quad (s_{j+1-n}, \dots, s_{j-1}, s_j) = (\bar{n}, \dots, \bar{2}, \bar{1}), \quad 1 \leq n \leq j,$$

obtained from (4.3) and (4.4), yields

$$\begin{aligned} \overline{H}_{2N-j+2}(z) &= \Delta(x^{2(N-j+1)}) \left(\sum_{n=1}^{j-1} \sum_{k=1}^{\text{Max}(j-n, n)} \sum_{\substack{\Omega_j \subset J_N \\ s_k=n \\ s_l=\bar{n}, l=j+1-n}} d_{\Omega_j \setminus \{s_l\}}(x, r) \overrightarrow{\Lambda}_{\Omega_j \setminus \{s_l\}}(x^{-1} z) \right. \\ &\quad \left. + \sum_{n=1}^j \sum_{k=1}^{\text{Max}(j+1-n, n)} \sum_{\substack{\Omega_j \subset J_N \\ s_{k-1} \prec n \prec s_k \\ s_l=\bar{n}, l=j+1-n}} d_{\Omega_j \setminus \{s_l\}}(x, r) \overrightarrow{\Lambda}_{\Omega_j \setminus \{s_l\}}(x^{-1} z) \right). \end{aligned}$$

Hence, we obtain $\overline{H}_{2N-j+2}(z) = \Delta(x^{2(N-j+1)}) T_{j-1}(x^{-1} z)$, $1 \leq j \leq N$.

We define $H_{2N-j+2}(z)$, $1 \leq j \leq N$, as the coefficient of $\delta(x^{2N-j+2} z_2/z_1)$ in (4.24). Adding the first term in (B.5), the second term in (B.5), the second term in (B.8), and the fourth term in (B.8) yields

$$\begin{aligned} H_{2N-j+2}(z) &= - \sum_{n=1}^{j-2} \sum_{l=n+2}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k=n, k=n+1 \\ s_l=\bar{n}}} d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2}) : \\ &\quad + \sum_{n=1}^{j-1} \sum_{l=n+1}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k=n, k=n \\ s_l=\bar{n}}} d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) : \\ &\quad + \sum_{n=1}^j \sum_{l=n}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k=n, k=n \\ s_l \prec \bar{n} \prec s_{l+1}}} \Delta(x^{2(N+1-l)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) : \\ &\quad - \sum_{n=1}^{j-1} \sum_{l=n+1}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k=n, k=n+1 \\ s_l \prec \bar{n} \prec s_{l+1}}} \Delta(x^{2(N+1-l)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) : . \end{aligned} \tag{4.28}$$

The first term in (4.28) vanishes, because there doesn't exist $s_1 \in J_N$, if $\Omega_j \subset J_N, s_k = n, k = n+1, s_l = \bar{n}, n+2 \leq l \leq j$, and $1 \leq n \leq j-2$ are satisfied. The fourth term in (4.28) vanishes, because there doesn't exist $s_1 \in J_N$, if $\Omega_j \subset J_N, s_k = n, k = n+1, s_l \prec \bar{n} \prec s_{l+1}, n+1 \leq l \leq j$, and $1 \leq n \leq j-1$ are satisfied. Rewriting the sum of the second and the third terms yields

$$\begin{aligned} H_{2N-j+2}(z) &= \sum_{n=1}^{j-1} \sum_{l=\text{Max}(n+1,j+1-n)}^j \sum_{\substack{\Omega_j \subset J_N \\ (s_1, s_2, \dots, s_n) = (1, 2, \dots, n) \\ s_l = \bar{n}}} d_{\Omega_j}(x, r) : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) : \\ &+ \sum_{n=1}^j \sum_{l=\text{Max}(n,j+1-n)}^j \sum_{\substack{\Omega_j \subset J_N \\ (s_1, s_2, \dots, s_n) = (1, 2, \dots, n) \\ s_l \prec \bar{n} \prec s_{l+1}}} \Delta(x^{2(N+1-l)}) d_{\Omega_j}(x, r) : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) : . \end{aligned}$$

The relation $d_{\Omega_j}(x, r) = \Delta(x^{2(N+1-j)}) d_{\Omega_j \setminus \{n\}}(x, r)$ holds, if $s_k = n, s_l = \bar{n}, 1 \leq n \leq j-1$, and $k = n$, and $n+1 \leq l \leq j$ are satisfied. The relation $d_{\Omega_j}(x, r) \Delta(x^{2(N+1-l)}) = \Delta(x^{2(N+1-j)}) d_{\Omega_j \setminus \{n\}}(x, r)$ holds, if $s_k = n, s_l \prec \bar{n} \prec s_{l+1}, 1 \leq n \leq j, k = n$, and $n+1 \leq l \leq j$ are satisfied. Using the above two relations and

$$: \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N-j+2} z) := \vec{\Lambda}_{\Omega_j \setminus \{n\}}(xz), \quad (s_1, s_2, \dots, s_n) = (1, 2, \dots, n), 1 \leq n \leq j,$$

obtained from (4.3) and (4.4), yields

$$\begin{aligned} H_{2N-j+2}(z) &= \Delta(x^{2(N-j+1)}) \left(\sum_{n=1}^{j-1} \sum_{l=\text{Max}(n+1,j+1-n)}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k = n, k = n \\ s_l = \bar{n}}} d_{\Omega_j \setminus \{s_k\}}(x, r) \vec{\Lambda}_{\Omega_j \setminus \{s_k\}}(xz) \right. \\ &\quad \left. + \sum_{n=1}^j \sum_{l=\text{Max}(n,j+1-n)}^j \sum_{\substack{\Omega_j \subset J_N \\ s_k = n, k = n \\ s_l \prec \bar{n} \prec s_{l+1}}} d_{\Omega_j \setminus \{s_k\}}(x, r) \vec{\Lambda}_{\Omega_j \setminus \{s_k\}}(xz) \right). \end{aligned}$$

Hence, we obtain $H_{2N-j+2}(z) = \Delta(x^{2(N-j+1)}) T_{j-1}(xz), 1 \leq j \leq N$.

We define $\bar{G}_{j+1-2m}(z), 1 \leq m \leq [\frac{j}{2}], 1 \leq j \leq N$, as the coefficient of $\delta(x^{-j-1+2m} z_2/z_1)$ in (4.24). Adding the first term in (B.2), the second term in (B.2), the first term in (B.7), the third term in (B.7), the first term in (B.8), and the third term in (B.8) yields

$$\bar{G}_{j+1-2m}(z) = \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_m \prec s \prec s_{m+1} \\ \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) :$$

$$\begin{aligned}
& - \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_{m-1} \prec s \prec s_m \\ \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\
& + \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{l=m+1 \\ s_m \prec s \prec s_{m+1} \\ s_l = \bar{s}, s=n}} \Delta(x^{2(l-m+n-N-1)}) d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\
& - \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{l=m \\ s_{m-1} \prec s \prec s_m \\ s_l = \bar{s}, s=n}} \Delta(x^{2(l-m+n-N-1)}) d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\
& + \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{k=1 \\ s_k = \bar{s}, s=\bar{n} \\ s_m \prec s \prec s_{m+1}}} \Delta(x^{2(m-k+n-N-1)}) d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\
& - \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{k=1 \\ s_k = \bar{s}, s=\bar{n} \\ s_{m-1} \prec s \prec s_m}} \Delta(x^{2(m-k+n-N-1)}) d_{\Omega_j}(x, r) : \Lambda_s(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_j}(z) : . \quad (4.29)
\end{aligned}$$

For a subset $\Omega_j = \{s_1, s_2, \dots, s_j\} \subset J_N$ and an element $s \notin \Omega_j, s \in J_N$, we write elements of $\Omega_j \cup \{s\}$ as $t_1, t_2, \dots, t_{j+1}, t_1 \prec t_2 \prec \dots \prec t_{j+1}$. In what follows, we use the abbreviation $\Omega_{j+1} = \{t_1, t_2, \dots, t_{j+1}\}$. Rewriting the sum yields

$$\begin{aligned}
& \bar{G}_{j+1-2m}(z) \\
& = \sum_{\substack{\Omega_{j+1} \subset J_N \\ \bar{t}_{m+1} \notin \Omega_{j+1} \setminus \{t_{m+1}\}}} d_{\{t_1, \dots, t_m\}}(x, r) d_{\{t_{m+2}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^m \prod_{\substack{q=m+2 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \quad \times : \Lambda_{t_{m+1}}(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_{j+1} \setminus \{t_{m+1}\}}(z) : \\
& \quad - \sum_{\substack{\Omega_{j+1} \subset J_N \\ \bar{t}_m \notin \Omega_{j+1} \setminus \{t_m\}}} d_{\{t_1, \dots, t_{m-1}\}}(x, r) d_{\{t_{m+1}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \quad \times : \Lambda_{t_m}(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_{j+1} \setminus \{t_m\}}(z) : \\
& \quad + \sum_{\Omega_{j+1} \subset J_N} \sum_{\substack{n=1 \\ t_{m+1}=n}}^N \sum_{\substack{l=m+2 \\ t_l=\bar{t}_{m+1}}}^{j+1} d_{\{t_1, \dots, t_m\}}(x, r) d_{\{t_{m+2}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^m \prod_{\substack{q=m+2 \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \quad \times \Delta(x^{2(l-m+t_{m+1}-N-2)}) : \Lambda_{t_{m+1}}(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_{j+1} \setminus \{t_{m+1}\}}(z) : \\
& \quad - \sum_{\Omega_{j+1} \subset J_N} \sum_{\substack{n=1 \\ t_m=n}}^N \sum_{\substack{l=m+1 \\ t_l=\bar{t}_m}}^{j+1} d_{\{t_1, \dots, t_{m-1}\}}(x, r) d_{\{t_{m+1}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \quad \times \Delta(x^{2(l-m+t_m-N-2)}) : \Lambda_{t_m}(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_{j+1} \setminus \{t_m\}}(z) : \\
& \quad + \sum_{\Omega_{j+1} \subset J_N} \sum_{\substack{n=1 \\ t_{m+1}=\bar{n}}}^N \sum_{\substack{k=1 \\ t_{m+1}=\bar{t}_k}}^m d_{\{t_1, \dots, t_m\}}(x, r) d_{\{t_{m+2}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^m \prod_{\substack{q=m+2 \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \quad \times \Delta(x^{2(m-k+t_k-N-1)}) : \Lambda_{t_{m+1}}(x^{-j-1+2m} z) \vec{\Lambda}_{\Omega_{j+1} \setminus \{t_{m+1}\}}(z) :
\end{aligned}$$

$$\begin{aligned}
& - \sum_{\Omega_{j+1} \subset J_N} \sum_{n=1}^N \sum_{\substack{k=1 \\ t_m = \bar{n} \\ t_m = \bar{t}_k}}^{m-1} d_{\{t_1, \dots, t_{m-1}\}}(x, r) d_{\{t_{m+1}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \Delta(x^{2(m-k+t_k-N-1)}) : \Lambda_{t_m}(x^{-j-1+2m} z) \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_m\}}(z) : .
\end{aligned}$$

Rewriting the sum yields

$$\begin{aligned}
& \overline{G}_{j+1-2m}(z) \\
& = \sum_{\Omega_{j+1} \subset J_N} d_{\{t_1, \dots, t_m\}}(x, r) d_{\{t_{m+2}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^m \prod_{\substack{q=m+2 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \prod_{\substack{p=1 \\ t_{m+1} = \bar{t}_p}}^m \Delta(x^{2(m-p+t_p-N-1)}) \prod_{\substack{q=m+2 \\ t_q = \bar{t}_{m+1}}}^{j+1} \Delta(x^{2(q-m+t_{m+1}-N-2)}) \times \\
& \times : \Lambda_{t_{m+1}}(x^{-j-1+2m} z) \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{m+1}\}}(z) : \\
& - \sum_{\Omega_{j+1} \subset J_N} d_{\{t_1, \dots, t_{m-1}\}}(x, r) d_{\{t_{m+1}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \prod_{\substack{p=1 \\ t_m = \bar{t}_p}}^m \Delta(x^{2(m-p+t_p-N-1)}) \prod_{\substack{q=m+1 \\ t_q = \bar{t}_m}}^{j+1} \Delta(x^{2(q-m+t_m-N-2)}) : \Lambda_{t_m}(x^{-j-1+2m} z) \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_m\}}(z) : .
\end{aligned}$$

Using

$$\begin{aligned}
& d_{\{t_1, \dots, t_m\}}(x, r) d_{\{t_{m+2}, \dots, t_{j+1}\}}(x, r) \times \\
& \times \prod_{p=1}^m \prod_{\substack{q=m+2 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \prod_{\substack{p=1 \\ t_{m+1} = \bar{t}_p}}^m \Delta(x^{2(m-p+t_p-N-1)}) \prod_{\substack{q=m+2 \\ t_q = \bar{t}_{m+1}}}^{j+1} \Delta(x^{2(q-m+t_{m+1}-N-2)}) \\
& = d_{\{t_1, \dots, t_{m-1}\}}(x, r) d_{\{t_{m+1}, \dots, t_{j+1}\}}(x, r) \times \\
& \times \prod_{p=1}^{m-1} \prod_{\substack{q=m+1 \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \prod_{\substack{p=1 \\ t_m = \bar{t}_p}}^m \Delta(x^{2(m-p+t_p-N-1)}) \prod_{\substack{q=m+1 \\ t_q = \bar{t}_m}}^{j+1} \Delta(x^{2(q-m+t_m-N-2)})
\end{aligned}$$

and

$$: \Lambda_{t_{m+1}}(x^{-j-1+2m} z) \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{m+1}\}}(z) :=: \Lambda_{t_m}(x^{-j-1+2m} z) \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_m\}}(z) :$$

yields $\overline{G}_{j+1-2m}(z) = 0$, $1 \leq m \leq [\frac{j}{2}]$, $1 \leq j \leq N$.

We define $G_{j+1-2m}(z)$, $1 \leq m \leq [\frac{j}{2}]$, $1 \leq j \leq N$, as the coefficient of $\delta(x^{j+1-2m} z_2/z_1)$ in (4.24). Adding the first term in (B.2), the second term in (B.2), the first term in (B.7), the

third term in (B.7), the first term in (B.8), and the third term in (B.8) yields

$$\begin{aligned}
& G_{j+1-2m}(z) \\
&= \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_{j+1-m} \prec s \prec s_{j+2-m} \\ \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : \\
&\quad - \sum_{\Omega_j \subset J_N} \sum_{\substack{s \in J_N \\ s_{j-m} \prec s \prec s_{j-m+1} \\ \bar{s} \notin \Omega_j}} d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : \\
&\quad + \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{l=j+2-m \\ s_{j-m+1} \prec s \prec s_{j-m+2} \\ s_l = \bar{s}, s=n}}^j \Delta(x^{2(l+m+n-j-N-2)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : \\
&\quad - \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{l=j+1-m \\ s_{j-m} \prec s \prec s_{j-m+1} \\ s_l = \bar{s}, s=n}}^j \Delta(x^{2(l+m+n-j-N-2)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : \\
&\quad + \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{k=1 \\ s_k = \bar{s}, s=\bar{n} \\ s_{j+1-m} \prec s \prec s_{j+2-m}}}^{j+1-m} \Delta(x^{2(j-m-k+n-N)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : \\
&\quad - \sum_{\Omega_j \subset J_N} \sum_{n=1}^N \sum_{\substack{k=1 \\ s_k = \bar{s}, s=\bar{n} \\ s_{j-m} \prec s \prec s_{j-m+1}}}^{j-m} \Delta(x^{2(j-m-k+n-N)}) d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_s(x^{j+1-2m} z) : .
\end{aligned} \tag{4.30}$$

Rewriting the sum, in the same way as the case of $\overline{G}_{j+1-2m}(z)$, yields

$$\begin{aligned}
& G_{j+1-2m}(z) \\
&= \sum_{\substack{\Omega_{j+1} \subset J_N \\ \bar{t}_{j+2-m} \notin \Omega_{j+1} \setminus \{t_{j+2-m}\}}} d_{\{t_1, \dots, t_{j+1-m}\}}(x, r) d_{\{t_{j+3-m}, \dots, t_{j+1}\}}(x, r) \times \\
&\quad \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=m+j+3-m \\ t_q = t_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+2-m}\}}(z) \Lambda_{t_{j+2-m}}(x^{j+1-2m} z) : \\
&\quad - \sum_{\substack{\Omega_{j+1} \subset J_N \\ \bar{t}_{j+1-m} \notin \Omega_{j+1} \setminus \{t_{j+1-m}\}}} d_{\{t_1, \dots, t_{j-m}\}}(x, r) d_{\{t_{j+2-m}, \dots, t_{j+1}\}}(x, r) \times \\
&\quad \times \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+1-m}\}}(z) \Lambda_{t_{j+1-m}}(x^{j+1-2m} z) : \\
&\quad + \sum_{\substack{\Omega_{j+1} \subset J_N \\ t_{j-m+2} = n}} \sum_{n=1}^N \sum_{\substack{l=j+2-m \\ t_l = \bar{n}}}^{j+1} d_{\{t_1, \dots, t_{j+1-m}\}}(x, r) d_{\{t_{j+3-m}, \dots, t_{j+1}\}}(x, r) \times \\
&\quad \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\ t_q = \bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \Delta(x^{2(l+m+n-j-N-2)}) \times
\end{aligned}$$

$$\begin{aligned}
& \times : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+2-m}\}}(z) \Lambda_{t_{j+2-m}}(x^{j+1-2m} z) : \\
& - \sum_{\Omega_{j+1} \subset J_N} \sum_{n=1}^N \sum_{\substack{l=j+1-m \\ t_l=\bar{n}}}^j d_{\{t_1, \dots, t_{j-m}\}}(x, r) d_{\{t_{j+2-m}, \dots, t_{j+1}\}}(x, r) \times \\
& \times \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \Delta(x^{2(l+m+n-j-N-2)}) : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+1-m}\}}(z) \Lambda_{t_{j+1-m}}(x^{j+1-2m} z) : \\
& + \sum_{\Omega_{j+1} \subset J_N} \sum_{\substack{n=1 \\ t_{j+2-m}=\bar{n}}}^N \sum_{\substack{k=1 \\ t_k=n}}^{j+1-m} d_{\{t_1, \dots, t_{j+1-m}\}}(x, r) d_{\{t_{j+3-m}, \dots, t_{j+1}\}}(x, r) \times \\
& \times \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \Delta(x^{2(j+n-m-k-N)}) : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+2-m}\}}(z) \Lambda_{t_{j+2-m}}(x^{j+1-2m} z) : \\
& - \sum_{\Omega_{j+1} \subset J_N} \sum_{\substack{n=1 \\ t_{j+1-m}=\bar{n}}}^N \sum_{\substack{k=1 \\ t_k=n}}^{j-m} d_{\{t_1, \dots, t_{j-m}\}}(x, r) d_{\{t_{j-m+1}, \dots, t_{j+1}\}}(x, r) \times \\
& \times \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \Delta(x^{2(j+n-m-k-N)}) : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+1-m}\}}(z) \Lambda_{t_{j+1-m}}(x^{j+1-2m} z) : .
\end{aligned}$$

Rewriting the sum yields

$$\begin{aligned}
& G_{j+1-2m}(z) \\
& = \sum_{\Omega_{j+1} \subset J_N} d_{\{t_1, \dots, t_{j+1-m}\}}(x, r) d_{\{t_{j+3-m}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \prod_{\substack{p=1 \\ t_{j+2-m}=\bar{t}_p}}^{j+1-m} \Delta(x^{2(j+t_p-m-p-N)}) \prod_{\substack{q=j+3-m \\ t_q=\bar{t}_{j+2-m}}}^{j+1} \Delta(x^{2(q+m+t_{j+2-m}-j-N-2)}) \times \\
& \times : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+2-m}\}}(z) \Lambda_{t_{j+2-m}}(x^{j+1-2m} z) : \\
& - \sum_{\Omega_{j+1} \subset J_N} d_{\{t_1, \dots, t_{j-m}\}}(x, r) d_{\{t_{j+2-m}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \prod_{\substack{p=1 \\ t_{j+1-m}=\bar{t}_p}}^{j-m} \Delta(x^{2(j+t_p-m-N-p)}) \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_{j+1-m}}}^{j+1} \Delta(x^{2(q+m+t_{j+1-m}-j-N-3)}) \times \\
& \times : \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+1-m}\}}(z) \Lambda_{t_{j+1-m}}(x^{j+1-2m} z) : .
\end{aligned}$$

Using

$$\begin{aligned}
& d_{\{t_1, \dots, t_{j+1-m}\}}(x, r) d_{\{t_{j+3-m}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{j+1-m} \prod_{\substack{q=j+3-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
& \times \prod_{\substack{p=1 \\ t_{j+2-m}=\bar{t}_p}}^{j+1-m} \Delta(x^{2(j+t_p-m-p-N)}) \prod_{\substack{q=j+3-m \\ t_q=\bar{t}_{j+2-m}}}^{j+1} \Delta(x^{2(q+m+t_{j+2-m}-j-N-3)}) = \\
&
\end{aligned}$$

$$\begin{aligned}
&= d_{\{t_1, \dots, t_{j-m}\}}(x, r) d_{\{t_{j-m+2}, \dots, t_{j+1}\}}(x, r) \prod_{p=1}^{j-m} \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_p}}^{j+1} \Delta(x^{2(q-p+t_p-N-2)}) \times \\
&\quad \times \prod_{\substack{p=1 \\ t_{j+1-m}=\bar{t}_p}}^{j-m} \Delta(x^{2(j+t_p-m-p-N)}) \prod_{\substack{q=j+2-m \\ t_q=\bar{t}_{j+1-m}}}^{j+1} \Delta(x^{2(q+m+t_{j+1-m}-j-N-3)})
\end{aligned}$$

and

$$: \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+2-m}\}}(z) \Lambda_{t_{j+2-m}}(x^{j+1-2m} z) :=: \overrightarrow{\Lambda}_{\Omega_{j+1} \setminus \{t_{j+1-m}\}}(z) \Lambda_{t_{j+1-m}}(x^{j+1-2m} z) :$$

yields $G_{j+1-2m}(z) = 0$, $1 \leq m \leq [\frac{j}{2}]$, $1 \leq j \leq N$.

We define $\overline{H}_{2N-j+2-2m}(z)$, $1 \leq j \leq N$, $1 \leq m \leq N - [\frac{j-1}{2}]$, as the coefficient of $\delta(x^{-2N+j-2+2m} z_2/z_1)$ in (4.24). We set

$$\overline{H}_{2N-j+2-2m}(z) = \sum_{\varepsilon=\pm} \varepsilon (\overline{\beta}_\varepsilon(z) + \overline{\gamma}_\varepsilon(z) + \overline{\delta}_\varepsilon(z)), \quad (4.31)$$

where we give $\overline{\beta}_+(z)$, $\overline{\beta}_-(z)$, $\overline{\gamma}_+(z)$, $\overline{\gamma}_-(z)$, $\overline{\delta}_+(z)$, and $\overline{\delta}_-(z)$ in (4.32), (4.33), (4.34), (4.35), (4.36), and (4.37), respectively. Adding the first term in (B.4) and the fourth term in (B.7) yields

$$\begin{aligned}
\overline{\beta}_+(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=m+1}^{\text{Min}(N, j+m-1)} \sum_{\substack{k=1 \\ s_k=n \\ s_l=\bar{n}, l=m+j+1-n}}^{j+m-n} d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j+2m-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\
&\quad + \sum_{\Omega_j \subset J_N} \sum_{n=m+1}^{\text{Min}(N, j+m)} \sum_{\substack{k=1 \\ s_{k-1} < n < s_k \\ s_l=\bar{n}, l=j+m+1-n}}^{j+m+1-n} d_{\Omega_j}(x, r) \Delta(x^{2(m+j-N-k)}) \times \\
&\quad \times : \Lambda_n(x^{-2N+j+2m-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \quad (4.32)
\end{aligned}$$

Adding the second term in (B.4) and the second term in (B.7) yields

$$\begin{aligned}
\overline{\beta}_-(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=m}^{\text{Min}(N, j+m-2)} \sum_{\substack{k=1 \\ s_k=n \\ s_l=\bar{n}, l=j+m-n}}^{j+m-n-1} d_{\Omega_j}(x, r) : \Lambda_n(x^{-2N+j+2m-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : \\
&\quad + \sum_{\Omega_j \subset J_N} \sum_{n=m}^{\text{Min}(N, j+m-1)} \sum_{\substack{k=1 \\ s_{k-1} < n < s_k \\ s_l=\bar{n}, l=j+m-n}}^{j+m-n} d_{\Omega_j}(x, r) \Delta(x^{2(m+j-N-k)}) \times \\
&\quad \times : \Lambda_n(x^{-2N+j+2m-2} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \quad (4.33)
\end{aligned}$$

Adding the first term in (B.6) yields

$$\overline{\gamma}_+(z) = \sum_{\substack{\Omega_j \subset J_N \\ s_k=0, k=j+m-N}} d_{\Omega_j}(x, r) : \Lambda_0(x^{-2N+j-2+2m} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \quad (4.34)$$

Adding the second term in (B.6) yields

$$\overline{\gamma}_-(z) = \sum_{\substack{\Omega_j \subset J_N \\ s_k=0, k=j+m-N-1}} d_{\Omega_j}(x, r) : \Lambda_0(x^{-2N+j-2+2m} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \quad (4.35)$$

Adding the first term in (B.5) and the fourth term in (B.8) yields

$$\begin{aligned} \bar{\delta}_+(z) = & \sum_{\Omega_j \subset J_N} \sum_{n=2N+2-j-m}^N \sum_{\substack{l=m+j+n-2N \\ s_k=n, k=j+m+n-2N-1 \\ s_l=\bar{n}}}^j d_{\Omega_j}(x, r) : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : \\ & + \sum_{\Omega_j \subset J_N} \sum_{n=2N+2-j-m}^N \sum_{\substack{l=m+j+n-2N-1 \\ s_k=n, k=j+m+n-2N-1 \\ s_l < \bar{n} < s_{l+1}}}^j d_{\Omega_j}(x, r) \Delta(x^{2(N+l+1-m-j)}) \times \\ & \times : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : . \end{aligned} \quad (4.36)$$

Adding the second term in (B.5) and the second term in (B.8) yields

$$\begin{aligned} \bar{\delta}_-(z) = & \sum_{\Omega_j \subset J_N} \sum_{n=2N+3-j-m}^N \sum_{\substack{l=j+m+n-2N-1 \\ s_k=n, k=j+m+n-2N-2 \\ s_l=\bar{n}}}^j d_{\Omega_j}(x, r) : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : \\ & + \sum_{\Omega_j \subset J_N} \sum_{n=2N+3-j-m}^N \sum_{\substack{l=j+m+n-2N-2 \\ s_k=n, k=j+m+n-2N-2 \\ s_l < \bar{n} < s_{l+1}}}^j d_{\Omega_j}(x, r) \Delta(x^{2(N+l+1-m-j)}) \times \\ & \times : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : . \end{aligned} \quad (4.37)$$

We show $\bar{H}_{2N-j+2-2m}(z) = 0, 1 \leq j \leq N, 1 \leq m \leq N - [\frac{j-1}{2}]$. In this proof we frequently use relation (4.4). The proof is divided into three cases: (i), $1 \leq m \leq N - j$, (ii), $m = N + 1 - j$, and (iii), $N + 2 - j \leq m \leq N - [\frac{j-1}{2}]$.

First, we study the case (i), $j + m \leq N$. In the case (i), $\bar{\gamma}_{\pm}(z)$ and $\bar{\delta}_{\pm}(z)$ vanish. Hence, we have $\bar{H}_{2N-j+2-2m}(z) = \bar{\beta}_+(z) - \bar{\beta}_-(z)$. We start from $\bar{\beta}_+(z)$. Rewriting the sum yields

$$\begin{aligned} \bar{\beta}_+(z) = & \left(\sum_{n=m+1}^{\text{Min}(N,j+m-1)} \sum_{k=1}^{j+m-n} \sum_{\substack{\Omega_j \subset J_N \\ s_k=n}} \sum_{\substack{(s_l, s_{l+1}, \dots, s_{l+r})=(\bar{n}, \bar{n-1}, \dots, \bar{n-r}) \\ n-r-1 < s_{l+r+1}, l=m+j+1-n}}^{n-m-1} \Delta(x^{2(m+j-N-k)}) \right. \\ & \left. + \sum_{n=m+1}^{\text{Min}(N,j+m)} \sum_{k=1}^{j+m-n-1} \sum_{\substack{\Omega_j \subset J_N \\ s_{k-1} < n < s_k}} \sum_{\substack{(s_l, s_{l+1}, \dots, s_{l+r})=(\bar{n}, \bar{n-1}, \dots, \bar{n-r}) \\ n-r-1 < s_{l+r+1}, l=m+j+1-n}}^{n-m-1} \Delta(x^{2(m+j-N-k)}) \right) \times \\ & \times \prod_{a=1}^{k-1} \prod_{\substack{b=1 \\ s_a=n-b}}^r \Delta(x^{2(m+j-N-a)}) \prod_{p=1}^{k-1} \prod_{\substack{q=l+r+1 \\ s_q=\bar{s}_p}}^j \Delta(x^{2(q-p+s_p-N-1)}) \times \\ & \times \prod_{\substack{k+1 \leq p < q \leq l-1 \\ s_q=\bar{s}_p}} \Delta(x^{2(q-p+s_p-N-1)}) : \Lambda_n(x^{-2N+j+2m-2} z) \vec{\Lambda}_{\Omega_j}(z) : . \end{aligned} \quad (4.38)$$

Using

$$\begin{aligned} & : \Lambda_n(x^{-2N+j+2m-2} z) \vec{\Lambda}_{\{s_1, s_2, \dots, s_{l-1}, \bar{n}, \bar{n-1}, \dots, \bar{n-r}, s_{l+r+1}, s_{l+r+2}, \dots, s_j\}}(z) : \\ & = : \Lambda_{n-r-1}(x^{-2N+j+2m-2} z) \vec{\Lambda}_{\{s_1, s_2, \dots, s_{l-1}, \bar{n-1}, \bar{n-2}, \dots, \bar{n-r-1}, s_{l+r+1}, s_{l+r+2}, \dots, s_j\}}(z) : , \\ & 0 \leq r \leq n - m - 1, l = m + j + 1 - n, \end{aligned} \quad (4.39)$$

obtained from (4.4), and replacing n by $n + 1$ yields $\bar{\beta}_+(z) = \bar{\beta}_-(z)$. Hence, we obtain $\bar{H}_{2N-j+2-2m}(z) = \bar{\beta}_+(z) - \bar{\beta}_-(z) = 0$ for $1 \leq m \leq N - j$.

Next, we examine the case (ii), $m = N + 1 - j$. In the case (ii), $\bar{\delta}_\pm(z)$ and $\bar{\gamma}_\pm(z)$ vanish. Hence, we have $\bar{H}_{2N-j+2-2m}(z) = \bar{\beta}_+(z) + \bar{\gamma}_+(z) - \bar{\beta}_-(z)$. We start from $\bar{\beta}_+(z) + \bar{\gamma}_+(z)$. Using (4.4) yields

$$\bar{\gamma}_+(z) = \Delta(1) \sum_{k=1}^j \sum_{\substack{\Omega_j \subset J_N \\ \overline{N+1-k} \prec s_{k+1} \prec \cdots \prec s_j}} : \Lambda_{N+1-k}(x^{-j} z) \vec{\Lambda}_{\{\overline{N}, \overline{N-1}, \dots, \overline{N+1-k}, s_{k+1}, s_{k+2}, \dots, s_j\}}(z) : . \quad (4.40)$$

Using (4.38), (4.39), and (4.40) yields $\bar{\beta}_+(z) + \bar{\gamma}_+(z) = \bar{\beta}_-(z)$. Hence, we obtain $\bar{H}_{2N-j+2-2m}(z) = \bar{\beta}_+(z) + \bar{\gamma}_+(z) - \bar{\beta}_-(z) = 0$ for $m = N - j + 1$.

Next, we examine the case (iii), $N + 2 - j \leq m \leq N - [\frac{j-1}{2}]$. Rewriting the sum yields

$$\begin{aligned} & \bar{\delta}_+(z) \\ &= \left(\sum_{n=2N+2-j-m}^N \sum_{l=j+m+n-2N}^j \sum_{\substack{\Omega_j \subset J_N \\ s_l = \bar{n} \\ (s_{k-r}, s_{k-r+1}, \dots, s_k) = (n-r, n-r+1, \dots, n) \\ s_{k-r-1} \prec n-r-1, k = m+j+n-2N-1}} \Delta(x^{2(l-j-m+N)}) \right. \\ &+ \sum_{n=2N+2-j-m}^N \sum_{l=m+j+n-2N-1}^j \sum_{\substack{\Omega_j \subset J_N \\ s_l = \bar{n} \\ (s_{k-r}, s_{k-r+1}, \dots, s_k) = (n-r, n-r+1, \dots, n) \\ s_{k-r-1} \prec n-r-1, k = m+j+n-2N-1}} \Delta(x^{2(l-j-m+N+1)}) \times \\ &\quad \times \prod_{a=1}^r \prod_{\substack{b=l+1 \\ \bar{s}_b = n-a}}^j \Delta(x^{2(N+b-j-m)}) \prod_{p=1}^{n-r-1} \prod_{\substack{q=l+1 \\ s_q = \bar{s}_p}}^j \Delta(x^{2(q-p+s_p-N-1)}) \times \\ &\quad \times \prod_{\substack{k+1 \leq p < q \leq l-1 \\ s_q = \bar{s}_p}} \Delta(x^{2(q-p+s_p-N-1)}) : \vec{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : \end{aligned} \quad (4.41)$$

and

$$\begin{aligned} \bar{\gamma}_+(z) &= \sum_{\substack{\Omega_j \subset J_N \\ s_k = 0, k = j+m-N \\ \overline{N} \prec s_{k+1}}} \prod_{p=1}^{k-1} \prod_{\substack{q=k+1 \\ s_q = \bar{s}_p}}^j \Delta(x^{2(q-p+s_p-N-1)}) : \vec{\Omega}_{\Lambda_j}(z) \Lambda_0(x^{-2N+j-2+2m}) : \\ &+ \sum_{\substack{\Omega_j \subset J_N \\ s_k = 0, k = j+m-N \\ \overline{N} \prec s_{k+1}}} \sum_{r=0}^{j-1} \prod_{a=1}^{k-1} \prod_{b=0}^r \Delta(x^{2(j+m-N-a)}) \times \\ &\quad \times \prod_{p=1}^{k-1} \prod_{\substack{q=k+r+2 \\ s_q = \bar{s}_p}}^j \Delta(x^{2(q-p+s_p-N-1)}) : \vec{\Omega}_{\Lambda_j}(z) \Lambda_0(x^{-2N+j-2+2m}) : . \end{aligned} \quad (4.42)$$

Using (4.2) and (4.4) yields

$$\begin{aligned} & : \vec{\Lambda}_{\{s_1, s_2, \dots, s_{k-r-1}, n-r, n-r+1, \dots, n, s_{k+1}, s_{k+2}, \dots, s_j\}}(z) \Lambda_{\bar{n}}(x^{-2N+j+2m-2} z) : \\ &= : \vec{\Lambda}_{\{s_1, s_2, \dots, s_{k-r-1}, n-r-1, n-r, \dots, n-1, s_{k+1}, s_{k+2}, \dots, s_j\}}(z) \Lambda_{\overline{n-r-1}}(x^{-2N+j+2m-2} z) : , \\ & 0 \leq r \leq k-1, k = j+m+n-2N-1, \end{aligned} \quad (4.43)$$

and

$$\begin{aligned} & : \overrightarrow{\Lambda}_{\{s_1, s_2, \dots, s_{k-1}, 0, \bar{N}, \bar{N-1}, \dots, \bar{N-r}, s_{k+r+2}, s_{k+r+3}, \dots, s_j\}}(z) \Lambda_0(x^{-2N+j+2m-2} z) : \\ & = \Delta(1) : \overrightarrow{\Lambda}_{\{s_1, s_2, \dots, s_{k-1}, \bar{N}, \bar{N-1}, \dots, \bar{N-r-1}, s_{k+r+2}, s_{k+r+3}, \dots, s_j\}}(z) \Lambda_{N-r-1}(x^{-2N+j+2m-2} z) : , \\ & \quad 0 \leq r \leq 2N + l + 1 - j - m, k = j + m + n - 2N - 1. \end{aligned} \quad (4.44)$$

Using (4.38), (4.39), (4.41), (4.42), (4.43), and (4.44) and replacing n by $n + 1$ yield $\bar{\beta}_+(z) + \bar{\gamma}_+(z) + \bar{\delta}_+(z) = \bar{\beta}_-(z) + \bar{\gamma}_-(z) + \bar{\delta}_-(z)$. Hence, we obtain $\bar{H}_{2N-j+2-2m}(z) = \sum_{\varepsilon=\pm} \varepsilon (\bar{\beta}_\varepsilon(z) + \bar{\gamma}_\varepsilon(z) + \bar{\delta}_\varepsilon(z)) = 0$ for $N + 2 - j \leq m \leq N - [\frac{j-1}{2}]$. Finally, we obtain $\bar{H}_{2N-j+2-2m}(z) = 0$ for $1 \leq m \leq N - [\frac{j-1}{2}]$.

We define $H_{2N-j+2-2m}(z)$, $1 \leq j \leq N$, $1 \leq m \leq N - [\frac{j-1}{2}]$, as the coefficient of $\delta(x^{-2N+j-2+2m} z_2/z_1)$ in (4.24). We set

$$H_{2N-j+2-2m}(z) = \begin{cases} \sum_{\varepsilon=\pm} \varepsilon (\beta_\varepsilon(z) + \gamma_\varepsilon(z) + \delta_\varepsilon(z)) & \text{otherwise,} \\ 0 & \text{if } j \text{ is even, } m = N - [\frac{j-1}{2}], \end{cases} \quad (4.45)$$

where we give $\beta_+(z)$, $\beta_-(z)$, $\gamma_+(z)$, $\gamma_-(z)$, $\delta_+(z)$, and $\delta_-(z)$ in (4.46), (4.47), (4.48), (4.49), (4.50), and (4.51), respectively. In (4.45) we define $H_0(z) = 0$ to avoid ambiguity of $\bar{H}_0(z)$ and $H_0(z)$. In the case when j is even, we have $\text{LHS}_{1,j} = c(x, r) (\bar{H}_0(z_2) - H_0(z_2)) \delta(z_2/z_1) + \bar{H}_2(z_2) \delta(x^{-2} z_2/z_1) - H_2(z_2) \delta(x^2 z_2/z_1) + \dots$. Adding the first term in (B.5) and the fourth term in (B.8) yields

$$\begin{aligned} \beta_+(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=m}^{\min(N, j+m-2)} \sum_{\substack{l=n+2-m \\ s_k=n, k=n+1-m \\ s_l=\bar{n}}}^j d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N+2-j-2m} z) : \\ &\quad + \sum_{\Omega_j \subset J_N} \sum_{n=m}^{\min(N, j+m-1)} \sum_{\substack{l=n+1-m \\ s_k=n, k=n+1-m \\ s_l \prec \bar{n} \prec s_{l+1}}}^j d_{\Omega_j}(x, r) \Delta(x^{2(l+m-N-1)}) \times \\ &\quad \times : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N+2-j-2m} z) : . \end{aligned} \quad (4.46)$$

Adding the second term in (B.5) and the second term in (B.8) yields

$$\begin{aligned} \beta_-(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=m+1}^{\min(N, j+m-1)} \sum_{\substack{l=n+1-m \\ s_k=n, k=n-m \\ s_l=\bar{n}}}^j d_{\Omega_j}(x, r) : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N+2-j-2m} z) : \\ &\quad + \sum_{\Omega_j \subset J_N} \sum_{n=m+1}^{\min(N, j+m)} \sum_{\substack{l=n-m \\ s_k=n, k=n-m \\ s_l \prec \bar{n} \prec s_{l+1}}}^j d_{\Omega_j}(x, r) \Delta(x^{2(l+m-N-1)}) \times \\ &\quad \times : \overrightarrow{\Lambda}_{\Omega_j}(z) \Lambda_{\bar{n}}(x^{2N+2-j-2m} z) : . \end{aligned} \quad (4.47)$$

Adding the first term in (B.6) yields

$$\gamma_+(z) = \sum_{\substack{\Omega_j \subset J_N \\ s_k=0, k=N+2-m}} d_{\Omega_j}(x, r) : \Lambda_0(x^{2N+2-j-2m} z) \overrightarrow{\Lambda}_{\Omega_j}(z) : . \quad (4.48)$$

Adding the second term in (B.6) yields

$$\gamma_-(z) = \sum_{\substack{\Omega_j \subset J_N \\ s_k=0, k=N+1-m}} d_{\Omega_j}(x, r) : \Lambda_0(x^{2N+2-j-2m} z) \vec{\Lambda}_{\Omega_j}(z) : . \quad (4.49)$$

Adding the first term in (B.4) and the fourth term in (B.7) yields

$$\begin{aligned} \delta_+(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=2N+3-j-m}^N \sum_{\substack{k=1 \\ s_k=n \\ s_l=\bar{n}, l=2N+3-m-n}}^{2N+2-m-n} d_{\Omega_j}(x, r) : \Lambda_n(x^{2N+2-j-2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\ &+ \sum_{\Omega_j \subset J_N} \sum_{n=2N+3-j-m}^N \sum_{\substack{k=1 \\ s_{k-1} < n < s_k \\ s_l=\bar{n}, l=2N+3-m-n}}^{2N+3-m-n} d_{\Omega_j}(x, r) \Delta(x^{2(N+2-m-k)}) \times \\ &\times : \Lambda_n(x^{2N+2-j-2m} z) \vec{\Lambda}_{\Omega_j}(z) : . \end{aligned} \quad (4.50)$$

Adding the second term in (B.4) and the second term in (B.7) yields

$$\begin{aligned} \delta_-(z) &= \sum_{\Omega_j \subset J_N} \sum_{n=2N+2-j-m}^N \sum_{\substack{k=1 \\ s_k=n \\ s_l=\bar{n}, l=2N+2-m-n}}^{2N+1-m-n} d_{\Omega_j}(x, r) : \Lambda_n(x^{2N+2-j-2m} z) \vec{\Lambda}_{\Omega_j}(z) : \\ &+ \sum_{\Omega_j \subset J_N} \sum_{n=2N+2-j-m}^N \sum_{\substack{k=1 \\ s_{k-1} < n < s_k \\ s_l=\bar{n}, l=2N+2-m-n}}^{2N+2-m-n} d_{\Omega_j}(x, r) \Delta(x^{2(N+2-m-k)}) \times \\ &\times : \Lambda_n(x^{2N+2-j-2m} z) \vec{\Lambda}_{\Omega_j}(z) : . \end{aligned} \quad (4.51)$$

The relation $H_{2N+2-j-2m}(z) = 0$ is shown in the same way as $\overline{H}_{2N+2-j-2m}(z) = 0$. ■

Lemma 4.7. *The currents $T_i(z)$ satisfy the following fusion relation*

$$\begin{aligned} &\lim_{z_1 \rightarrow x^{\pm(i+j)} z_2} \left(1 - \frac{x^{\pm(i+j)} z_2}{z_1} \right) f_{i,j} \left(\frac{z_2}{z_1} \right) T_i(z_1) T_j(z_2) \\ &= \mp c(x, r) \prod_{l=1}^{\text{Min}(i,j)-1} \Delta(x^{2l+1}) T_{i+j}(x^{\pm i} z_2), \quad 1 \leq i, j \leq N. \end{aligned} \quad (4.52)$$

Proof. For subsets $\Omega_i^{(1)} = \{s_1, s_2, \dots, s_i\} \subset J_N$ with $s_1 \prec s_2 \prec \dots \prec s_i$ and $\Omega_j^{(2)} = \{t_1, t_2, \dots, t_j\} \subset J_N$ with $t_1 \prec t_2 \prec \dots \prec t_j$, we set $\Omega_{i+j} = \Omega_i^{(1)} \cup \Omega_j^{(2)}$. From (4.1), the necessary and sufficient condition that $f_{i,j}(z_2/z_1) \vec{\Lambda}_{\Omega_i^{(1)}}(z_1) \vec{\Lambda}_{\Omega_j^{(2)}}(z_2)$ has a pole at $z_1 = x^{-(i+j)} z_2$ (respectively $z_1 = x^{i+j} z_2$) is $s_i \prec t_1$ (respectively $t_j \prec s_1$). In the case when $s_i \prec t_1$ or $t_j \prec s_1$, we obtain

$$\begin{aligned} f_{i,j} \left(\frac{z_2}{z_1} \right) \vec{\Lambda}_{\Omega_i^{(1)}}(z_1) \vec{\Lambda}_{\Omega_j^{(2)}}(z_2) &= \prod_{k=0}^{\text{Min}(i,j)-1} \Delta \left(\frac{x^{\pm(2k+1-i-j)} z_2}{z_1} \right) \times \\ &\times \prod_{p=1}^i \prod_{\substack{q=1 \\ t_q=\bar{s}_p}}^j \Delta \left(\frac{x^{\pm\{2(q-p+s_p-N-1+i)-i-j\}} z_2}{z_1} \right) : \vec{\Lambda}_{\Omega_i^{(1)}}(z_1) \vec{\Lambda}_{\Omega_j^{(2)}}(z_2) : . \end{aligned}$$

The signs \pm in the products in the above expression of $f_{i,j}(z_2/z_1) \overrightarrow{\Lambda}_{\Omega_i^{(1)}}(z_1) \overrightarrow{\Lambda}_{\Omega_j^{(2)}}(z_2)$ are in the same order. The upper sign is for $s_i \prec t_1$, and the lower sign is for $t_j \prec s_1$. Taking the limit yields

$$\begin{aligned} & \lim_{z_1 \rightarrow x^{\pm(i+j)} z_2} \left(1 - \frac{x^{\pm(i+j)} z_2}{z_1} \right) f_{i,j} \left(\frac{z_2}{z_1} \right) \overrightarrow{\Lambda}_{\Omega_i^{(1)}}(z_1) \overrightarrow{\Lambda}_{\Omega_j^{(2)}}(z_2) \\ &= \mp c(x, r) \prod_{l=1}^{\text{Min}(i,j)-1} \Delta(x^{2l+1}) \prod_{p=1}^i \prod_{\substack{q=1 \\ t_q = \bar{s}_p}}^j \Delta \left(x^{2(q-p+s_p-N-1+i)} \right) \overrightarrow{\Lambda}_{\Omega_{i+j}}(x^{\pm i} z_2), \quad 1 \leq i, j \leq N. \end{aligned} \quad (4.53)$$

Here, we use : $\overrightarrow{\Lambda}_{\Omega_i^{(1)}}(x^{\pm(i+j)} z) \overrightarrow{\Lambda}_{\Omega_j^{(2)}}(z) := \overrightarrow{\Lambda}_{\Omega_{i+j}}(x^{\pm i} z)$. Adding (4.53) over all $\Omega_i^{(1)}$ and $\Omega_j^{(2)}$ yields (4.52). \blacksquare

Lemma 4.8. *The currents $T_i(z)$ satisfy the following fusion relations*

$$\begin{aligned} & \lim_{z_1 \rightarrow x^{\pm(2N+1+i-j)} z_2} \left(1 - \frac{x^{\pm(2N+1+i-j)} z_2}{z_1} \right) f_{i,j} \left(\frac{z_2}{z_1} \right) T_i(z_1) T_j(z_2) \\ &= \mp c(x, r) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) \prod_{l=N+1-j}^{N+i-j} \Delta(x^{2l}) T_{j-i}(x^{\pm i} z_2), \quad 1 \leq i \leq j \leq N, \end{aligned} \quad (4.54)$$

$$\begin{aligned} & \lim_{z_1 \rightarrow x^{\pm(2N+1-i+j)} z_2} \left(1 - \frac{x^{\pm(2N+1-i+j)} z_2}{z_1} \right) f_{i,j} \left(\frac{z_2}{z_1} \right) T_i(z_1) T_j(z_2) \\ &= \mp c(x, r) \prod_{l=1}^{j-1} \Delta(x^{2l+1}) \prod_{l=N+1-i}^{N+j-i} \Delta(x^{2l}) T_{i-j}(x^{\pm(2N+1-i)} z_2), \quad 1 \leq j \leq i \leq N. \end{aligned} \quad (4.55)$$

Proof. Using (3.4), (4.8), (4.9), and (4.52) yields (4.54) and (4.55). \blacksquare

Proof. Here we will give a proof of Theorem 3.2. We prove Theorem 3.2 by induction. Lemma 4.6 is the base for induction. We define $\text{LHS}_{i,j}$, $\text{RHS1}_{i,j}$ and $\text{RHS2}_{i,j}(k)$ with $1 \leq k \leq i \leq j \leq N$ as

$$\begin{aligned} \text{LHS}_{i,j} &= f_{i,j} \left(\frac{z_2}{z_1} \right) T_i(z_1) T_j(z_2) - f_{j,i} \left(\frac{z_1}{z_2} \right) T_j(z_2) T_i(z_1), \\ \text{RHS1}_{i,j} &= c(x, r) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) \prod_{l=N+1-j}^{N+i-j} \Delta(x^{2l}) \times \\ &\quad \times \left(\delta \left(\frac{x^{-2N+j-i-1} z_2}{z_1} \right) T_{j-i}(x^{-i} z_2) - \delta \left(\frac{x^{2N-j+i+1} z_2}{z_1} \right) T_{j-i}(x^i z_2) \right), \\ \text{RHS2}_{i,j}(k) &= c(x, r) \prod_{l=1}^{k-1} \Delta(x^{2l+1}) \left(\delta \left(\frac{x^{-j+i-2k} z_2}{z_1} \right) f_{i-k,j+k}(x^{j-i}) T_{i-k}(x^k z_1) T_{j+k}(x^{-k} z_2) \right. \\ &\quad \left. - \delta \left(\frac{x^{j-i+2k} z_2}{z_1} \right) f_{i-k,j+k}(x^{-j+i}) T_{i-k}(x^{-k} z_1) T_{j+k}(x^k z_2) \right), \quad 1 \leq k \leq i-1, \\ \text{RHS2}_{i,j}(i) &= c(x, r) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) \left(\delta \left(\frac{x^{-j-i} z_2}{z_1} \right) T_{j+i}(x^{-i} z_2) - \delta \left(\frac{x^{j+i} z_2}{z_1} \right) T_{j+i}(x^i z_2) \right). \end{aligned}$$

We prove the following relation by induction on $i, 1 \leq i \leq j \leq N$.

$$\text{LHS}_{i,j} = \text{RHS1}_{i,j} + \sum_{k=1}^i \text{RHS2}_{i,j}(k). \quad (4.56)$$

The base, $i = 1 \leq j \leq N$ was proved previously in Lemma 4.6.

We assume that the relation (4.56) holds for some $i, 1 \leq i < j \leq N$, and show that $\text{LHS}_{i+1,j} = \text{RHS1}_{i+1,j} + \sum_{k=1}^{i+1} \text{RHS2}_{i+1,j}(k)$ from this assumption. First, we summarize some relations. The assumption (4.56) yields

$$\lim_{w_1 \rightarrow x^{2N-i-j} w_2} \left(1 - x^{2N-i-j} \frac{w_2}{w_1} \right) f_{j-1,i} \left(\frac{w_2}{w_1} \right) T_{j-1}(w_1) T_i(w_2) = 0, \quad (4.57)$$

$$\lim_{w_1 \rightarrow x^{2N-i-j} w_2} \left(1 - x^{2N-i-j} \frac{w_2}{w_1} \right) f_{1,j-i} \left(\frac{w_2}{w_1} \right) T_1(w_1) T_{j-i}(w_2) = 0, \quad (4.58)$$

$$f_{i,j} \left(\frac{w_2}{w_1} \right) T_i(w_1) T_j(w_2) = f_{j,i} \left(\frac{w_1}{w_2} \right) T_j(w_2) T_i(w_1) \quad (4.59)$$

for $\frac{w_2}{w_1} \neq x^{\pm(2N-j+i+1)}, x^{\pm(j-i+2k)}, 1 \leq k \leq i$. Direct calculation yields

$$\lim_{w_2 \rightarrow x^{-i-1} w_1} \left(1 - x^{-i-1} \frac{w_1}{w_2} \right) \Delta \left(x^{-i} \frac{w_1}{w_2} \right) = c(x, r). \quad (4.60)$$

Multiplying $\text{LHS}_{i,j}$ by $f_{1,i}(z_1/z_3) f_{1,j}(z_2/z_3) T_1(z_3)$ on the left and using the quadratic relation (4.56) with $i = 1$, along with the fusion relation (4.10) yields

$$\begin{aligned} & f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{LHS}_{i,j} \\ &= f_{1,j} \left(\frac{z_2}{z_3} \right) f_{i,j} \left(\frac{z_2}{z_1} \right) f_{1,i} \left(\frac{z_1}{z_3} \right) T_1(z_3) T_i(z_1) T_j(z_2) \\ & \quad - f_{j,1} \left(\frac{z_3}{z_2} \right) f_{j,i} \left(\frac{z_1}{z_2} \right) T_j(z_2) f_{1,i} \left(\frac{z_1}{z_3} \right) T_1(z_3) T_i(z_1) \\ & \quad - c(x, r) \Delta(x^{2(N+1-j)}) \delta \left(\frac{x^{-2N+j-2} z_2}{z_3} \right) \Delta \left(\frac{x^{-i} z_1}{z_3} \right) f_{j-1,i} \left(\frac{x^{-2N+j-1} z_1}{z_3} \right) \times \\ & \quad \times T_{j-1}(x^{2N-j+1} z_3) T_i(z_1) \\ & \quad + c(x, r) \Delta(x^{2(N+1-j)}) \delta \left(\frac{x^{2N-j+2} z_2}{z_3} \right) \Delta \left(\frac{x^i z_1}{z_3} \right) f_{j-1,i} \left(\frac{x^{2N-j+1} z_1}{z_3} \right) \times \\ & \quad \times T_{j-1}(x^{-2N+j-1} z_3) T_i(z_1) \\ & \quad - c(x, r) \delta \left(\frac{x^{-j-1} z_2}{z_3} \right) \Delta \left(\frac{x^{-i} z_1}{z_3} \right) f_{j+1,i} \left(\frac{x^{-j} z_1}{z_3} \right) T_{j+1}(x^j z_3) T_i(z_1) \\ & \quad + c(x, r) \delta \left(\frac{x^{j+1} z_2}{z_3} \right) \Delta \left(\frac{x^i z_1}{z_3} \right) f_{j+1,i} \left(\frac{x^j z_1}{z_3} \right) T_{j+1}(x^{-j} z_3) T_i(z_1). \end{aligned} \quad (4.61)$$

Taking the limit $z_3 \rightarrow x^{-i-1} z_1$ of (4.61) multiplied by $c(x, r)^{-1} (1 - x^{-i-1} z_1/z_3)$ and using the

relations (4.52) and (4.55), (4.57), (4.59), and (4.60) yields

$$\begin{aligned}
& \lim_{z_3 \rightarrow x^{-i-1} z_1} \frac{1}{c(x, r)} \left(1 - x^{-i-1} \frac{z_1}{z_3} \right) f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{LHS}_{i,j} \\
& = f_{i+1,j} \left(\frac{x z_2}{z_1} \right) T_{i+1}(x^{-1} z_1) T_j(z_2) - f_{j,i+1} \left(\frac{x^{-1} z_1}{z_2} \right) T_j(z_2) T_{i+1}(x^{-1} z_1) \\
& + c(x, r) \prod_{l=1}^i \Delta(x^{2l+1}) \prod_{l=N-j+1}^{N+i+1-j} \Delta(x^{2l}) \delta \left(\frac{x^{2N-j+i+3} z_2}{z_1} \right) T_{j-i-1}(x^{i+1} z_2) \\
& - c(x, r) \delta \left(\frac{x^{i-j} z_2}{z_1} \right) f_{i,j+1}(x^{-i+j+1}) T_i(z_1) T_{j+1}(x^{-1} z_2) \\
& + c(x, r) \delta \left(\frac{x^{i+j+2} z_2}{z_1} \right) \prod_{l=1}^i \Delta(x^{2l+1}) T_{i+j+1}(x^{i+1} z_2). \tag{4.62}
\end{aligned}$$

Multiplying $\text{RHS1}_{i,j}$ by $f_{1,i}(z_1/z_3) f_{1,j}(z_2/z_3) T_1(z_3)$ from the left and using fusion relations (4.9) and (4.10) yields

$$\begin{aligned}
& f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS1}_{i,j} \\
& = c(x, r) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) \prod_{l=N+1-j}^{N+i-j} \Delta(x^{2l}) \times \left\{ \delta \left(\frac{x^{-2N+j-i-1} z_2}{z_1} \right) \Delta \left(\frac{x^i z_1}{z_3} \right) f_{1,j-i} \left(\frac{x^{-i} z_2}{z_3} \right) \times \right. \\
& \quad \left. \times T_1(z_3) T_{j-i}(x^{-i} z_2) - \delta \left(\frac{x^{2N-j+i+1} z_2}{z_1} \right) \Delta \left(\frac{x^{-i} z_1}{z_3} \right) f_{1,j-i} \left(\frac{x^i z_2}{z_3} \right) T_1(z_3) T_{j-i}(x^i z_2) \right\}. \tag{4.63}
\end{aligned}$$

Taking the limit $z_3 \rightarrow x^{-i-1} z_1$ of (4.63) multiplied by $c(x, r)^{-1} (1 - x^{-i-1} z_1/z_3)$ and using the relations (4.54) and (4.58) yields

$$\begin{aligned}
& \lim_{z_3 \rightarrow x^{-i-1} z_1} \frac{1}{c(x, r)} \left(1 - x^{-i-1} \frac{z_1}{z_3} \right) f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS1}_{i,j} \\
& = c(x, r) \prod_{l=1}^i \Delta(x^{2l+1}) \prod_{l=N+1-j}^{N+i+1-j} \Delta(x^{2l}) \delta \left(\frac{x^{-2N+j-i-1} z_2}{z_1} \right) T_{j-i-1}(x^{-i-1} z_2). \tag{4.64}
\end{aligned}$$

Multiplying $\text{RHS2}_{i,j}(i)$ by $f_{1,i}(z_1/z_3) f_{1,j}(z_2/z_3) T_1(z_3)$ from the left and using the fusion relation (4.11) yields

$$\begin{aligned}
& f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS2}_{i,j}(i) \\
& = c(x, r) \prod_{l=1}^{i-1} \Delta_1(x^{2l+1}) \left(\delta \left(\frac{x^{-i-j} z_2}{z_1} \right) f_{1,i+1} \left(\frac{x^j z_1}{z_3} \right) \Delta \left(\frac{x^i z_1}{z_3} \right) T_1(z_3) T_{i+j}(x^j z_1) \right. \\
& \quad \left. - \delta \left(\frac{x^{i+j} z_2}{z_1} \right) f_{1,i+1} \left(\frac{x^{-j} z_1}{z_3} \right) \Delta \left(\frac{x^{-i} z_1}{z_3} \right) T_1(z_3) T_{i+j}(x^{-j} z_1) \right). \tag{4.65}
\end{aligned}$$

Taking the limit $z_3 \rightarrow x^{-i-1} z_1$ of (4.65) multiplied by $c(x, r)^{-1} (1 - x^{-i-1} z_1/z_3)$ and using the

relations (4.52) and (4.60) yields

$$\begin{aligned}
& \lim_{z_3 \rightarrow x^{-i-1} z_1} \frac{1}{c(x, r)} \left(1 - x^{-i-1} \frac{z_1}{z_3} \right) f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS2}_{i,j}(i) \\
& = c(x, r) \delta \left(\frac{x^{-i-j} z_2}{z_1} \right) \prod_{l=1}^i \Delta(x^{2l+1}) T_{i+j+1}(x^{-i-1} z_2) \\
& \quad - c(x, r) \delta \left(\frac{x^{i+j} z_2}{z_1} \right) \prod_{l=1}^{i-1} \Delta(x^{2l+1}) f_{1,i+j}(x^{i-j+1}) T_1(x^{-i-1} z_1) T_{i+j}(x^i z_2). \tag{4.66}
\end{aligned}$$

Multiplying $\text{RHS2}_{i,j}(k)$, $1 \leq k \leq i-1$, by $f_{1,i}(z_1/z_3) f_{1,j}(z_2/z_3) T_1(z_3)$ from the left and using relations (4.12) and (4.59) yields

$$\begin{aligned}
& f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS2}_{i,j}(k) \\
& = c(x, r) \prod_{l=1}^{k-1} \Delta(x^{2l+1}) \left(\delta \left(\frac{x^{-j+i-2k} z_2}{z_1} \right) f_{1,i-k} \left(\frac{x^k z_1}{z_3} \right) f_{j+k,i-k}(x^{i-j}) f_{1,j+k} \left(\frac{x^{-i+j+k} z_1}{z_3} \right) \right. \\
& \quad \times T_1(z_3) T_{j+k}(x^{j-i+k} z_1) T_{i-k}(x^k z_1) \\
& \quad - \delta \left(\frac{x^{j-i+2k} z_2}{z_1} \right) f_{1,i-k} \left(\frac{x^{-k} z_1}{z_3} \right) f_{i-k,j+k}(x^{i-j}) f_{1,j+k} \left(\frac{x^{i-j-k} z_1}{z_3} \right) \\
& \quad \left. \times T_1(z_3) T_{i-k}(x^{-k} z_1) T_{j+k}(x^k z_2) \right), 1 \leq k \leq i-1. \tag{4.67}
\end{aligned}$$

Taking the limit $z_3 \rightarrow x^{-i-1} z_1$ of (4.67) multiplied by $c(x, r)^{-1} (1 - x^{-i-1} z_1/z_3)$, and using the fusion relations (4.7), (4.52), and (4.59) yields

$$\begin{aligned}
& \lim_{z_3 \rightarrow x^{-i-1} z_1} \frac{1}{c(x, r)} \left(1 - x^{-i-1} \frac{z_1}{z_3} \right) f_{1,i} \left(\frac{z_1}{z_3} \right) f_{1,j} \left(\frac{z_2}{z_3} \right) T_1(z_3) \times \text{RHS2}_{i,j}(k) \\
& = c(x, r) \prod_{l=1}^k \Delta(x^{2l+1}) \delta \left(\frac{x^{-j+i-2k} z_2}{z_1} \right) f_{j+k-1,i-k}(x^{i-j+1}) T_{i-k}(x^k z_1) T_{j+k+1}(x^{-k-1} z_2) \\
& \quad - c(x, r) \prod_{l=1}^{k-1} \Delta(x^{2l+1}) \delta \left(\frac{x^{j-i+2k} z_2}{z_1} \right) f_{i-k+1,j+k}(x^{i-j+1}) T_{i-k+1}(x^{-k-1} z_1) T_{j+k}(x^k z_2), \\
& \quad 1 \leq k \leq i-1. \tag{4.68}
\end{aligned}$$

Adding (4.62), (4.64), (4.66), and (4.68) for $1 \leq k \leq i-1$, and replacing z_1 by $x z_1$ yields $\text{LHS}_{i+1,j} = \text{RHS1}_{i+1,j} + \sum_{k=1}^{i+1} \text{RHS2}_{i+1,j}(k)$. By induction on i , we proved quadratic relation (3.5). \blacksquare

4.3 Proof of Lemma 3.4

Lemma 4.9. *The current $T_1(z)$ commutes with the screening currents $S_k(w)$ as follows.*

$$[T_1(z), S_k(w)] = C_k(z) (D_{x^r} \delta) \left(\frac{x^k w}{z} \right) + \bar{C}_k(z) (D_{x^r} \delta) \left(\frac{x^{2N+1-k} w}{z} \right), \quad 1 \leq k \leq N. \tag{4.69}$$

Here we set q -difference

$$(D_q \delta)(z) = \delta(qz) - \delta(q^{-1}z),$$

the currents $C_k(z)$ and $\bar{C}_k(z)$, $1 \leq k \leq N$, are given by

$$\begin{aligned} C_k(z) &= x^{-r+1}(x^{r-1} - x^{-r+1}) : \Lambda_k(z) S_k(x^{r-k} z) :, \\ \bar{C}_k(z) &= x^{r-1}(x^{r-1} - x^{-r+1}) : \Lambda_{\bar{k}}(z) S_k(x^{-2N-1+k-r} z) : . \end{aligned}$$

Proof. Adding (B.10) yields

$$\begin{aligned} &[T_1(z), S_k(w)] \\ &= (x^{r-1} - x^{-r+1}) \left(-x^{-r+1} : \Lambda_k(z) S_k(w) : \delta \left(x^{k-r} \frac{w}{z} \right) + x^{r-1} : \Lambda_{k+1}(z) S_k(w) : \delta \left(x^{k+r} \frac{w}{z} \right) \right. \\ &\quad \left. + x^{r-1} : \Lambda_{\bar{k}}(z) S_k(w) : \delta \left(x^{2N+1-k+r} \frac{w}{z} \right) - x^{-r+1} : \Lambda_{\bar{k+1}}(z) S_k(w) : \delta \left(x^{2N+1-k-r} \frac{w}{z} \right) \right), \\ &1 \leq k \leq N-1, \end{aligned}$$

and

$$\begin{aligned} &[T_1(z), S_N(w)] \\ &= (x^{-r+1} - x^{r-1}) \left(x^{-r+1} : \Lambda_N(z) S_N(w) : \delta \left(x^{N-r} \frac{w}{z} \right) - x^{r-1} : \Lambda_{\bar{N}}(z) S_{\bar{N}}(w) : \delta \left(x^{N+1+r} \frac{w}{z} \right) \right) \\ &\quad + \frac{[r-1]_x [\frac{1}{2}]_x}{[r-\frac{1}{2}]_x} (x - x^{-1}) \left(\delta \left(x^{N+r} \frac{w}{z} \right) - \delta \left(x^{N+1-r} \frac{w}{z} \right) \right) : \Lambda_0(z) S_N(w) : . \end{aligned}$$

Using the relations

$$\begin{aligned} x^{-r+1} : \Lambda_k(z) S_k(x^{r-k} z) &:= x^{r-1} : \Lambda_{k+1}(z) S_k(x^{-r-k} z) :, \quad 1 \leq k \leq N-1, \\ x^{-r+1} : \Lambda_N(z) S_N(x^{r-N} z) &:= \frac{[\frac{1}{2}]_x}{[r-\frac{1}{2}]_x} : \Lambda_0(z) S_N(x^{-r-N} z) :, \\ x^{r-1} : \Lambda_{\bar{k}}(z) S_k(x^{-2N-1+k-r} z) &:= x^{-r+1} : \Lambda_{\bar{k+1}}(z) S_k(x^{-2N-1+k+r} z) :, \quad 1 \leq k \leq N-1, \\ x^{r-1} : \Lambda_{\bar{N}}(z) S_N(x^{-r-N-1} z) &:= \frac{[\frac{1}{2}]_x}{[r-\frac{1}{2}]_x} : \Lambda_0(z) S_N(x^{r-N-1} z) : . \end{aligned}$$

yields (4.69). ■

Corollary 4.10. The current $T_1(z)$ commutes with the screening operators S_k

$$[T_1(z), S_k] = 0, \quad 1 \leq k \leq N. \quad (4.70)$$

Proof. From (4.69), we obtain

$$[T_1(z), S_k] = \oint \frac{dw}{2\pi\sqrt{-1}w} \left(C_k(z) (D_{x^r} \delta) \left(\frac{x^k w}{z} \right) + \bar{C}_k(z) (D_{x^r} \delta) \left(\frac{x^{2N+1-k} w}{z} \right) \right).$$

Using $\oint \frac{dw}{2\pi\sqrt{-1}w} (D_{x^r} \delta) \left(\frac{x^s w}{z} \right) = 0$ with $s = k, 2N+1-k$ yields $[T_1(z), S_k] = 0$. ■

Proof. Here we will give a proof of Lemma 3.4. Set $T_j(z) = \sum_{m \in \mathbf{Z}} T_j[m] z^{-m}$, $1 \leq j \leq 2N$ and

$$f_{i,j}(z) = \sum_{l=0}^{\infty} f_{i,j}^l z^l. \text{ From (4.21), we obtain}$$

$$\begin{aligned} &(x^{-(j+1)k+m} - x^{(j+1)k-m}) T_{j+1}[m] \\ &= \Delta(x^{2N+2-2j}) (x^{(2N-j+2)k-m} - x^{(-2N+j-2)k+m}) T_{j-1}[m] \\ &\quad + c(x, r)^{-1} \sum_{l=0}^{\infty} \left(f_{1,j}^l T_1[k-l] T_j[l-k+m] - f_{j,1}^l T_j[k-l-m] T_1[l-k] \right), \quad m, k \in \mathbf{Z}, 1 \leq j \leq N. \end{aligned}$$

Hence, $T_{j+1}[m], m \in \mathbf{Z}, 1 \leq j \leq N$, are expressed in terms of $T_j[n], T_{j-1}[n]$, and $T_1[n], n \in \mathbf{Z}, 1 \leq j \leq N$. From duality (3.4), $T_j[m], m \in \mathbf{Z}, N+2 \leq j \leq 2N$ are expressed in terms of $T_{2N+1-j}[n], n \in \mathbf{Z}, N+2 \leq j \leq 2N$. Finally, $T_j[m], m \in \mathbf{Z}, 1 \leq j \leq 2N$ are expressed in terms of $T_1[n], n \in \mathbf{Z}$. Hence, we obtain (3.6) from (4.70). \blacksquare

5 Conclusion and discussion

In this paper, we obtained the free field construction of higher W -currents $T_i(z), i \geq 2$, of the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$. We obtained a closed set of quadratic relations for the W -currents $T_i(z)$, which are completely different from those in types $A_N^{(1)}$ and $A(M,N)^{(1)}$. The quadratic relations of $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ do not preserve “parity”, though those of $\mathcal{W}_{x,r}(A_N^{(1)})$ and $\mathcal{W}_{x,r}(A(M,N)^{(1)})$ do. Here we define “parity” of $T_i(z)T_j(w)$ as $i+j$. We obtained the duality $T_{2N+1-i}(z) = c_i T_i(z), 1 \leq i \leq N$, which is a new structure that does not occur in types $A_2^{(2)}$, $A_N^{(1)}$, and $A(M,N)^{(1)}$. This allowed us to define the deformed W -algebra $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ using generators and relations similarly to the definition of the twisted affine Lie algebra of type $A_{2N}^{(2)}$ given in Section 2.

We also justified our definition of the deformed W -algebra of type $A_{2N}^{(2)}$. We compare Definition 3.3 with other definitions. In Ref.[6], the deformed W -algebras of types $A_N^{(1)}, B_N^{(1)}, C_N^{(1)}, D_N^{(1)}$, and $A_{2N}^{(2)}$ were proposed as the intersection of the kernels of the screening operators. We recall the definition based on the screening operators for $A_{2N}^{(2)}$. Let $\mathbf{H}_{x,r}$ be the vector space spanned by the formal power series currents of the form

$$: \partial_z^{n_1} Y_{i_1}(x^{rj_1+k_1} z)^{\varepsilon_1} \cdots \partial_z^{n_l} Y_{i_l}(x^{rj_l+k_l} z)^{\varepsilon_l} :,$$

where $\varepsilon_i = \pm 1$ ²⁾. We define $\mathbf{W}_{x,r}$ as the vector subspace of $\mathbf{H}_{x,r}$ consisting of all currents that commute with the screening operators $S_i, 1 \leq i \leq N$, in (2.7). Let $\{F_a(z) = \sum_{m \in \mathbf{Z}} F_a[m] z^{-m}\}_{a \in A}$ be a basis of the vector space $\mathbf{W}_{x,r}$. Let \mathcal{W}^{FR} be the associative algebra generated by elements $F_a[m], m \in \mathbf{Z}, a \in A$. Let J_K be the left ideal of \mathcal{W}^{FR} generated by elements $F_a[m], m \geq K \in \mathbf{N}, a \in A$. We define the deformed W -algebra

$$\mathcal{W}_{x,r}^{FR}(A_{2N}^{(2)}) = \varprojlim \mathcal{W}^{FR} / J_K.$$

We propose another definition of the deformed W -algebra. From (3.6), the W -currents $T_i(z) = \sum_{m \in \mathbf{Z}} T_i[m] z^{-m}, 1 \leq i \leq 2N$, commute with the screening operators. Let \mathcal{W}^{AKOS} be the associative algebra generated by elements $T_i[m], m \in \mathbf{Z}, 1 \leq i \leq 2N$. Let L_K be the left ideal of \mathcal{W}^{AKOS} generated by elements $T_i[m], m \geq K \in \mathbf{N}, 1 \leq i \leq 2N$. We define the deformed W -algebra

$$\mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)}) = \varprojlim \mathcal{W}^{AKOS} / L_K.$$

In this study, our definitions $\mathcal{W}_{x,r}(A_{2N}^{(2)})$ were based on generators and relations. We have introduced three definitions of the deformed W -algebra for the twisted algebra of the type $A_{2N}^{(2)}$.

²⁾ We define $Y_i(z)^{-1}$ as the inverse element of $Y_i(z)$, that is, $Y_i(z)Y_i(z)^{-1} = Y_i(z)^{-1}Y_i(z) = 1$. Specifically, we obtain $Y_i(z)^{-1} = x^{-ry_i(0)} \langle Y_i(z)Y_i(z) \rangle : \exp \left(-\sum_{m \neq 0} y_i(m) z^{-m} \right) :$, where we used the symbol $\langle \rangle$ defined in (A.1).

Conjecture 1. $\mathcal{W}_{x,r}(A_{2N}^{(2)})$, $\mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)})$, and $\mathcal{W}_{x,r}^{FR}(A_{2N}^{(2)})$ are isomorphic as associative algebras.

$$\mathcal{W}_{x,r}(A_{2N}^{(2)}) \cong \mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)}) \cong \mathcal{W}_{x,r}^{FR}(A_{2N}^{(2)}). \quad (5.1)$$

The author believes that this conjecture can be extended to arbitrary affine Lie algebras. Some necessary conditions of isomorphism (5.1) in Conjecture 1 can be indicated immediately. From (3.6), we obtain the following inclusion:

$$\mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)}) \subseteq \mathcal{W}_{x,r}^{FR}(A_{2N}^{(2)}).$$

We establish a homomorphism of associative algebras $\varphi \in \text{Hom}_{\mathbf{C}}(\mathcal{W}_{x,r}(A_{2N}^{(2)}), \mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)}))$ using $\varphi(\bar{T}_i[m]) = T_i[m]$. φ is surjective,

$$\varphi(\mathcal{W}_{x,r}(A_{2N}^{(2)})) = \mathcal{W}_{x,r}^{AKOS}(A_{2N}^{(2)}).$$

If we assume that φ is injective, the isomorphism on the left side in (5.1) is obtained. In other words, no independent relations other than (3.4) and (3.5) exist in $\mathcal{W}_{x,r}(A_{2N}^{(2)})$. We propose two results to support this claim. In the classical limit the second Hamiltonian structure $\{\cdot, \cdot\}$ of the q -Poisson algebra [6, 17, 18, 19] was obtained from the quadratic relations (see (3.7) and (3.8)). In the conformal limit all defining relations of the W -algebra $\mathcal{W}_{\beta}(A_N^{(1)})$, $N = 1, 2$, are obtained from the quadratic relations of $\mathcal{W}_{x,r}(A_N^{(1)})$ upon the assumption that the currents $T_i(z)$ have the form of expansion for small parameter \hbar (see Appendix of Ref.[3]).

The definition of the deformed W -algebra $\mathcal{W}_{x,r}(\mathfrak{g})$ for non-twisted affine Lie algebra \mathfrak{g} was formulated in terms of the quantum Drinfeld–Sokolov reduction in Ref.[9]. Formulating the definition of the deformed W -algebras $\mathcal{W}_{x,r}(\mathfrak{g})$ in terms of the quantum Drinfeld–Sokolov reduction for twisted affine Lie algebra or affine Lie superalgebra [11, 12, 13, 14, 15] is still a problem that needs to be solved.

It remains an open challenge to identify quadratic relations of the deformed W -algebras $\mathcal{W}_{x,r}(\mathfrak{g})$ for the affine Lie algebras \mathfrak{g} except for types $A_N^{(1)}$ and $A_{2N}^{(2)}$. We believe that this paper presents a key step towards extending our construction for general affine Lie algebras \mathfrak{g} . In [6] and [11] the free field construction of the basic W -current $T_1(z)$ of $\mathcal{W}_{x,r}(\mathfrak{g})$ was suggested in the case when the underlying simple finite-dimensional Lie algebra $\overset{\circ}{\mathfrak{g}}$ is of classical type,

$$T_1(z) = \begin{cases} \Lambda_1(z) + \cdots + \Lambda_{N+1}(z) & \text{for } \mathfrak{g} \text{ of type } A_N^{(1)}, \\ \Lambda_1(z) + \cdots + \Lambda_N(z) + \Lambda_0(z) + \Lambda_{\bar{N}}(z) + \cdots + \Lambda_{\bar{1}}(z) & \text{for } \mathfrak{g} \text{ of types } B_N^{(1)}, A_{2N}^{(2)}, D_{N+1}^{(2)}, \\ \Lambda_1(z) + \cdots + \Lambda_N(z) + \Lambda_{\bar{N}}(z) + \cdots + \Lambda_{\bar{1}}(z) & \text{for } \mathfrak{g} \text{ of types } C_N^{(1)}, D_N^{(1)}, A_{2N-1}^{(2)}. \end{cases}$$

Here we omit details of free field constructions of $\Lambda_i(z)$. The free field construction of $T_1(z)$ has similar form to that for \mathfrak{g} of type $A_{2N}^{(2)}$ except for the case of $A_N^{(1)}$. Therefore, we expect that a similar duality as (3.4) and similar quadratic relation (3.5) hold in all cases in types $B_N^{(1)}, C_N^{(1)}, D_N^{(1)}, A_{2N-1}^{(2)}$, and $D_{N+1}^{(2)}$. We would like to draw your attention to the following analogy. Let \mathfrak{g} be an affine Lie algebras of one of the types $B_N^{(1)}, C_N^{(1)}, D_N^{(1)}, A_{2N-1}^{(2)}$, or $D_{N+1}^{(2)}$. Let $\overset{\circ}{\mathfrak{g}}$ be the underlying simple finite-dimensional Lie algebra. Let $\overset{\circ}{\mathfrak{h}}$ be a Cartan subalgebra of $\overset{\circ}{\mathfrak{g}}$. Let $\bar{\Lambda}_1, \bar{\Lambda}_2, \dots, \bar{\Lambda}_l$ be the fundamental weights of $\overset{\circ}{\mathfrak{g}}$, where l is the dimension of $\overset{\circ}{\mathfrak{h}}$. Let $V_{\bar{\Lambda}_1}$ be the integrable highest weight representation of $U_q(\overset{\circ}{\mathfrak{g}})$ with the highest weight $\bar{\Lambda}_1$. Let V be the evaluation representation corresponding to $V_{\bar{\Lambda}_1}$ of the quantum affine algebra $U_q(\mathfrak{g})$ with a spectral parameter $z \in \mathbf{C}^\times$. Let n be the dimension of $V_{\bar{\Lambda}_1}$. We have $\overset{n-i}{\wedge} V \simeq \left(\overset{i}{\wedge} V\right)^* \simeq \overset{i}{\wedge} V^*$,

because $\wedge^n V \simeq \mathbf{C}$. The evaluation representation V of $U_q(\mathfrak{g})$ is self-dual except for \mathfrak{g} of type $A_N^{(1)}$. Hence, we obtain the duality of the representations of $U_q(\mathfrak{g})$,

$$\wedge^{n-i} V \xrightarrow{i} \wedge V \text{ if } \mathfrak{g} \text{ is not of type } A_N^{(1)},$$

which is similar as that in (3.4). As an analogy, we expect the duality of the W -currents,

$$T_{n-i}(z) = c_i T_i(z) \text{ if } \mathfrak{g} \text{ is not of type } A_N^{(1)},$$

for the deformed W algebras $\mathcal{W}_{x,r}(\mathfrak{g})$. Here $c_i, 0 \leq i \leq n$, are constants.

It remains an open challenge to identify quadratic relations of the deformed W -algebras $\mathcal{W}_{x,r}(\mathfrak{g})$ for affine superalgebra \mathfrak{g} except for those of type $A(M,N)^{(1)}$. Recently the deformed W -superalgebra $\mathcal{W}_{x,r}(\mathfrak{g})$ has appeared in the study of D-branes and physical interest is growing to this subject, see e.g. [15]. As revealed in Refs. [13, 14, 15], it is expected that, in cases of superalgebras \mathfrak{g} , infinite number of higher W -currents $T_i(z), i = 1, 2, 3, \dots$, satisfy a closed set of infinite number of quadratic relations. It is interesting to understand how duality will be extended to the case of superalgebras. We expect to report on quadratic relations and duality for more general deformed W -algebras $\mathcal{W}_{x,r}(\mathfrak{g})$ associated with affine Lie algebras and affine Lie superalgebras in the near future.

A Normal ordering rules

We list the normal ordering rules. For operators $V(z)$ and $W(w)$ we use the notation

$$V(z)W(w) = \langle V(z)W(w) \rangle : V(z)W(w) : \quad (\text{A.1})$$

and write down only the part $\langle V(z)W(w) \rangle$ in the formulas below. Using the standard formula

$$e^A e^B = e^{[A,B]} e^B e^A \quad ([A, B], A] = 0 \text{ and } [[A, B], B] = 0),$$

we obtain the normal ordering rules.

A.1 $A_i(z)$ and $S_i(z)$

$$\begin{aligned} \langle A_i(z_1)A_i(z_2) \rangle &= \left(\Delta\left(\frac{xz_2}{z_1}\right) \Delta\left(\frac{x^{-1}z_2}{z_1}\right) \right)^{-1}, \quad 1 \leq i \leq N-1, \\ \langle A_N(z_1)A_N(z_2) \rangle &= \Delta\left(\frac{z_2}{z_1}\right) \left(\Delta\left(\frac{xz_2}{z_1}\right) \Delta\left(\frac{x^{-1}z_2}{z_1}\right) \right)^{-1}, \\ \langle A_i(z_1)A_j(z_2) \rangle &= \Delta\left(\frac{z_2}{z_1}\right), \quad |i-j|=1, 1 \leq i, j \leq N, \\ \langle A_i(z_1)A_j(z_2) \rangle &= 1, \quad |i-j| \geq 2, 1 \leq i, j \leq N, \end{aligned} \quad (\text{A.2})$$

$$\begin{aligned} \langle S_i(z_1)S_i(z_2) \rangle &= z_1^{\frac{2(r-1)}{r}} \left(1 - \frac{z_2}{z_1}\right) \frac{(x^2 z_2/z_1; x^{2r})_\infty}{(x^{2r-2} z_2/z_1; x^{2r})_\infty}, \quad 1 \leq i \leq N-1, \\ \langle S_N(z_1)S_N(z_2) \rangle &= z_1^{\frac{r-1}{r}} \left(1 - \frac{z_2}{z_1}\right) \frac{(x^2 z_2/z_1; x^{2r})_\infty (x^{2r-2} z_2/z_1; x^{2r})_\infty}{(x z_2/z_1; x^{2r})_\infty (x^{2r-1} z_2/z_1; x^{2r})_\infty}, \\ \langle S_i(z_1)S_j(z_2) \rangle &= z_1^{-\frac{r-1}{r}} \frac{(x^{2r-1} z_2/z_1; x^{2r})_\infty}{(x z_2/z_1; x^{2r})_\infty}, \quad |i-j|=1, 1 \leq i, j \leq N, \\ \langle S_i(z_1)S_j(z_2) \rangle &= 1, \quad |i-j| \geq 2, 1 \leq i, j \leq N, \end{aligned} \quad (\text{A.3})$$

$$\begin{aligned}
\langle A_i(z_1)S_i(z_2) \rangle &= x^{-4(r-1)} \frac{\left(1 - x^r \frac{z_2}{z_1}\right) \left(1 - x^{r-2} \frac{z_2}{z_1}\right)}{\left(1 - x^{-r} \frac{z_2}{z_1}\right) \left(1 - x^{2-r} \frac{z_2}{z_1}\right)}, \quad 1 \leq i \leq N-1, \\
\langle A_N(z_1)S_N(z_2) \rangle &= x^{-2(r-1)} \frac{\left(1 - x^r \frac{z_2}{z_1}\right) \left(1 - x^{r-2} \frac{z_2}{z_1}\right) \left(1 - x^{1-r} \frac{z_2}{z_1}\right)}{\left(1 - x^{-r} \frac{z_2}{z_1}\right) \left(1 - x^{2-r} \frac{z_2}{z_1}\right) \left(1 - x^{r-1} \frac{z_2}{z_1}\right)}, \\
\langle A_i(z_1)S_j(z_2) \rangle &= x^{2(r-1)} \frac{\left(1 - x^{1-r} \frac{z_2}{z_1}\right)}{\left(1 - x^{r-1} \frac{z_2}{z_1}\right)}, \quad |i-j|=1, 1 \leq i, j \leq N, \\
\langle A_i(z_1)S_j(z_2) \rangle &= 1, \quad |i-j| \geq 2, 1 \leq i, j \leq N,
\end{aligned} \tag{A.4}$$

$$\begin{aligned}
\langle S_i(z_1)A_i(z_2) \rangle &= \frac{\left(1 - x^{-r} \frac{z_2}{z_1}\right) \left(1 - x^{2-r} \frac{z_2}{z_1}\right)}{\left(1 - x^r \frac{z_2}{z_1}\right) \left(1 - x^{r-2} \frac{z_2}{z_1}\right)}, \quad 1 \leq i \leq N-1, \\
\langle S_N(z_1)A_N(z_2) \rangle &= \frac{\left(1 - x^{-r} \frac{z_2}{z_1}\right) \left(1 - x^{2-r} \frac{z_2}{z_1}\right) \left(1 - x^{r-1} \frac{z_2}{z_1}\right)}{\left(1 - x^r \frac{z_2}{z_1}\right) \left(1 - x^{r-2} \frac{z_2}{z_1}\right) \left(1 - x^{1-r} \frac{z_2}{z_1}\right)}, \\
\langle S_j(z_1)A_i(z_2) \rangle &= \frac{\left(1 - x^{r-1} \frac{z_2}{z_1}\right)}{\left(1 - x^{1-r} \frac{z_2}{z_1}\right)}, \quad |i-j|=1, 1 \leq i, j \leq N, \\
\langle S_j(z_1)A_i(z_2) \rangle &= 1, \quad |i-j| \geq 2, 1 \leq i, j \leq N.
\end{aligned} \tag{A.5}$$

A.2 $Y_i(z)$, $A_i(z)$ and $S_i(z)$

The symmetric matrix $I(m) = (I_{i,j}(m))_{i,j=1}^N$ is the inverse matrix of $B(m)$. The elements $I_{i,j}(m) = I_{j,i}(m)$, $1 \leq i \leq j \leq N$, are written as

$$I_{i,j}(m) = \frac{1}{[(N+1)m]_x - [Nm]_x} \times \begin{cases} [(N+1-j)m]_x - [(N-j)m]_x, & i = 1, 1 \leq j \leq N, \\ (-1)^{N-j+i} \sum_{k=i-1}^{N-j+i} (-1)^k [km]_x, & 2 \leq i \leq j \leq N-1, \\ [im]_x, & 1 \leq i \leq N, j = N. \end{cases} \tag{A.6}$$

The generators $y_i(m)$, $1 \leq i \leq N$, are written as

$$y_i(m) = \sum_{j=1}^N I_{i,j}(m) a_j(m), \quad Q_i^y = \sum_{j=1}^N I_{i,j}(0) Q_j. \tag{A.7}$$

From (2.2), (2.3) and (A.6) we obtain

$$\langle Y_1(z_1)Y_1(z_2) \rangle = f_{1,1} \left(\frac{z_2}{z_1} \right)^{-1},$$

$$\begin{aligned}
\langle Y_1(z_1)A_1(z_2) \rangle &= \Delta \left(\frac{z_2}{z_1} \right)^{-1}, \quad \langle Y_1(z_1)A_i(z_2) \rangle = 1, \quad 2 \leq i \leq N, \\
\langle A_1(z_1)Y_1(z_2) \rangle &= \Delta \left(\frac{z_2}{z_1} \right)^{-1}, \quad \langle A_i(z_1)Y_1(z_2) \rangle = 1, \quad 2 \leq i \leq N, \\
\langle Y_1(z_1)S_1(z_2) \rangle &= x^{-2(r-1)} \frac{\left(1 - x^{r-1} \frac{z_2}{z_1}\right)}{\left(1 - x^{1-r} \frac{z_2}{z_1}\right)}, \quad \langle Y_1(z_1)S_i(z_2) \rangle = 1, \quad 2 \leq i \leq N, \\
\langle S_1(z_1)Y_1(z_2) \rangle &= \frac{\left(1 - x^{1-r} \frac{z_2}{z_1}\right)}{\left(1 - x^{r-1} \frac{z_2}{z_1}\right)}, \quad \langle S_i(z_1)Y_1(z_2) \rangle = 1, \quad 2 \leq i \leq N.
\end{aligned} \tag{A.8}$$

B Exchange relations

In this appendix we list the exchange relations.

B.1 $\Lambda_i(z)$

We give the exchange relations of $\Lambda_j(z)$ and $\vec{\Lambda}_{\Omega_i}(z)$, which are obtained from (4.1). We set $s \in J_N = \{1, 2, \dots, N, 0, \bar{N}, \dots, \bar{2}, \bar{1}\}$. For an element $s \in J_N$ and a subset $\Omega_i = \{s_1, s_2, \dots, s_i\} \subset J_N$ with $s_1 \prec s_2 \prec \dots \prec s_i$, we calculate

$$X_{\Omega_i, s}(z_1, z_2) = f_{1,i} \left(\frac{z_2}{z_1} \right) \Lambda_s(z_1) \vec{\Lambda}_{\Omega_i}(z_2) - f_{i,1} \left(\frac{z_1}{z_2} \right) \vec{\Lambda}_{\Omega_i}(z_2) \Lambda_s(z_1). \tag{B.1}$$

- In the case of $s, \bar{s} \notin \Omega_i$, we obtain

$$X_{\Omega_i, s}(z_1, z_2) = c(x, r) : \Lambda_s(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{-i-3+2k} z_2}{z_1} \right) - \delta \left(\frac{x^{-i-1+2k} z_2}{z_1} \right) \right). \tag{B.2}$$

Here we set $k, 1 \leq k \leq i+1$, by $k = \begin{cases} 1 & \text{if } s \prec s_1, \\ q & \text{if } s_{q-1} \prec s \prec s_q, 2 \leq q \leq i, \\ i+1 & \text{if } s_i \prec s. \end{cases}$

- In the case of $s \in \Omega_i$ and $\bar{s} \notin \Omega_i$, we obtain

$$f_{1,i} \left(\frac{z_2}{z_1} \right) \Lambda_s(z_1) \vec{\Lambda}_{\Omega_i}(z_2) - f_{i,1} \left(\frac{z_1}{z_2} \right) \vec{\Lambda}_{\Omega_i}(z_2) \Lambda_s(z_1) = 0. \tag{B.3}$$

- In the case of $s, \bar{s} \in \Omega_i$ and $s = n, 1 \leq n \leq N$, we obtain

$$X_{\Omega_i, s}(z_1, z_2) = c(x, r) : \Lambda_n(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{-2N-i+2n+2l-4} z_2}{z_1} \right) - \delta \left(\frac{x^{-2N-i+2n+2l-2} z_2}{z_1} \right) \right). \tag{B.4}$$

Here we set $k, l, 1 \leq k < l \leq i$, by $s = n = s_k$ and $\bar{s} = \bar{n} = s_l$.

- In the case of $s, \bar{s} \in \Omega_i$ and $s = \bar{n}, 1 \leq n \leq N$, we obtain

$$X_{\Omega_i, s}(z_1, z_2) = c(x, r) : \Lambda_{\bar{n}}(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{2N-i-2n+2k} z_2}{z_1} \right) - \delta \left(\frac{x^{2N-i-2n+2k+2} z_2}{z_1} \right) \right). \tag{B.5}$$

Here we set $k, l, 1 \leq k < l \leq i$, by $\bar{s} = n = s_k$ and $s = \bar{n} = s_l$.

- In the case of $s = 0 \in \Omega_i$, we obtain

$$X_{\Omega_i, s}(z_1, z_2) = c(x, r) : \Lambda_0(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{-i-2+2k} z_2}{z_1} \right) - \delta \left(\frac{x^{-i+2k} z_2}{z_1} \right) \right). \tag{B.6}$$

Here we set $k, 1 \leq k \leq i$, by $s_k = 0$.

- In the case of $s \notin \Omega_i$ and $\bar{s} \in \Omega_i$ and $s = n, 1 \leq n \leq N$, we obtain

$$\begin{aligned} & X_{\Omega_i, s}(z_1, z_2) \\ &= c(x, r) \Delta(x^{2(l-k+n-N)}) : \Lambda_n(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{-i+2k-3} z_2}{z_1} \right) - \delta \left(\frac{x^{-2N+2n+2l-i-2} z_2}{z_1} \right) \right) \\ &+ c(x, r) \Delta(x^{2(l-k+n-N-1)}) : \Lambda_n(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(-\delta \left(\frac{x^{-i+2k-1} z_2}{z_1} \right) + \delta \left(\frac{x^{-2N+2n+2l-i-4} z_2}{z_1} \right) \right). \end{aligned} \quad (\text{B.7})$$

Here we set $k, l, 1 \leq k \leq l \leq i$, by $s_l = \bar{s} = \bar{n}$ and $k = \begin{cases} 1 & \text{if } s = n \prec s_1, \\ q & \text{if } s_{q-1} \prec s = n \prec s_q, 2 \leq q \leq i. \end{cases}$

- In the case of $s \notin \Omega_i$ and $\bar{s} \in \Omega_i$ and $s = \bar{n}, 1 \leq n \leq N$, we obtain

$$\begin{aligned} & X_{\Omega_i, s}(z_1, z_2) \\ &= c(x, r) \Delta(x^{2(l-k+n-N-1)}) : \Lambda_{\bar{n}}(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(\delta \left(\frac{x^{-i+2l-1} z_2}{z_1} \right) - \delta \left(\frac{x^{2N-2n-i+2k+2} z_2}{z_1} \right) \right) \\ &+ c(x, r) \Delta(x^{2(l-k+n-N)}) : \Lambda_{\bar{n}}(z_1) \vec{\Lambda}_{\Omega_i}(z_2) : \left(-\delta \left(\frac{x^{-i+2l+1} z_2}{z_1} \right) + \delta \left(\frac{x^{2N-2n-i+2k} z_2}{z_1} \right) \right). \end{aligned} \quad (\text{B.8})$$

Here we set $k, l, 1 \leq k \leq l \leq i$, by $s_k = \bar{s} = n$ and $l = \begin{cases} q & \text{if } s_q \prec s = \bar{n} \prec s_{q+1}, 1 \leq q \leq i-1, \\ i & \text{if } s_i \prec s = \bar{n}. \end{cases}$

B.2 $S_i(z)$

From (A.3) we obtain

$$\begin{aligned} S_i(z_1) S_i(z_2) &= -\frac{[u_2 - u_1 + 1]}{[u_1 - u_2 + 1]} S_i(z_2) S_i(z_1), \quad 1 \leq i \leq N-1, \\ S_N(z_1) S_N(z_2) &= -\frac{[u_1 - u_2 + \frac{1}{2}][u_2 - u_1 + 1]}{[u_2 - u_1 + \frac{1}{2}][u_1 - u_2 + 1]} S_N(z_2) S_N(z_1), \\ S_i(z_1) S_j(z_2) &= \frac{[u_1 - u_2 + \frac{1}{2}]}{[u_2 - u_1 + \frac{1}{2}]} S_j(z_2) S_i(z_1), \quad |i-j|=1, 1 \leq i, j \leq N, \\ S_i(z_1) S_j(z_2) &= S_j(z_2) S_i(z_1), \quad |i-j| \geq 2, 1 \leq i, j \leq N. \end{aligned} \quad (\text{B.9})$$

Here we set $z_i = x^{2u_i}$ ($i = 1, 2$) and $[u] = x^{\frac{u^2}{r}-2u} \Theta_{x^{2r}}(z)$.

B.3 $\Lambda_i(z)$ and $S_i(z)$

From (A.4), (A.5) and (A.8) we obtain

$$\begin{aligned} [\Lambda_k(z_1), S_k(z_2)] &= (x^{-2r+2} - 1) : \Lambda_k(z_1) S_k(z_2) : \delta \left(\frac{x^{k-r} z_2}{z_1} \right), \quad 1 \leq k \leq N, \\ [\Lambda_{k+1}(z_1), S_k(z_2)] &= (x^{2r-2} - 1) : \Lambda_{k+1}(z_1) S_k(z_2) : \delta \left(\frac{x^{k+r} z_2}{z_1} \right), \quad 1 \leq k \leq N-1, \\ [\Lambda_{\bar{k}}(z_1), S_k(z_2)] &= (x^{2r-2} - 1) : \Lambda_{\bar{k}}(z_1) S_k(z_2) : \delta \left(\frac{x^{2N+1-k+r} z_2}{z_1} \right), \quad 1 \leq k \leq N, \end{aligned} \quad (\text{B.10})$$

$$[\Lambda_{\overline{k+1}}(z_1), S_k(z_2)] = (x^{-2r+2} - 1) : \Lambda_{\overline{k+1}}(z_1) S_k(z_2) : \delta \left(\frac{x^{2N+1-k-r} z_2}{z_1} \right), \quad 1 \leq k \leq N-1,$$

$$[\Lambda_0(z_1), S_N(z_2)] = (x - x^{-1}) \frac{[r-1]_x [\frac{1}{2}]_x}{[r-\frac{1}{2}]_x} \left(\delta \left(\frac{x^{r+N} z_2}{z_1} \right) - \delta \left(\frac{x^{-r+N+1} z_2}{z_1} \right) \right) : \Lambda_0(z_1) S_N(z_2) : .$$

Other commutators on the type $[\Lambda_i(z_1), S_k(z_2)]$ that are used in the proof of Lemma 4.9 are zeroes.

Dedication

This paper is dedicated to Professor Michio Jimbo on the occasion of his 70th anniversary.

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